



Simulation of the Cryogenic Behavior
of the Superconducting Magnets
in the Energy Saver

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Abstract

A time-dependent model of the Energy Saver cryoloop has been developed to simulate the cryogenic behavior of the superconducting magnets. Preliminary results include temperature, pressure, enthalpy, and velocity distributions. These calculations are reported for steady-state conditions, and transient behavior during accelerator ramping and ramp-off. Transient effects on beam storage operation require further work.

Introduction

A Satellite refrigerator provides cooling to two parallel 20-magnet cryoloops. A single cryoloop is modeled with one-half Satellite flow; Fig. 1 shows the relationship of these parts of the system. Fig. 2 diagrams the 20-magnet string modeled here. Liquid helium flows to the magnet string through a subcooler at 1.8 atm, and is expanded at the far end to a lower pressure to return as two-phase liquid. The ramping heat load is deposited in the single-phase stream and is transferred to the colder two-phase counter-flow stream. The heat is primarily absorbed by the heat of vaporization of the two-phase fluid that returns to the Satellite through the shell side of the subcooler.

Model

The cryoloop was simulated by simultaneous numerical solution of the one-dimensional mass, energy, and momentum conservation equations describing the helium streams, subject to appropriate boundary and initial conditions. The method follows that used by M. C. Jones¹ on the counterflow cooldown of superconducting power transmission lines, and uses the computer program PDECOL developed by N. K. Madsen and R. F. Sincovec² to integrate the partial differential equations.

PDECOL solves a set of three equations for each stream: the one-dimensional conservation equations written in terms of pressure P , enthalpy H , and velocity V .³

$$\frac{\partial P}{\partial t} = -V \frac{\partial P}{\partial X} - c^2 \frac{\partial V}{\partial X} + \phi \left[\frac{\Lambda}{a} + \rho \frac{V^2 f |V| p}{2a} \right]$$

$$\frac{\partial H}{\partial t} = -V \frac{\partial H}{\partial X} - c^2 \frac{\partial V}{\partial X} + \left(\frac{1+\phi}{\rho} \right) \left[\frac{\Lambda}{a} + \rho \frac{V^2 f |V| p}{2a} \right]$$

$$\frac{\partial V}{\partial t} = -\frac{1}{\rho} \frac{\partial P}{\partial X} - V \frac{\partial V}{\partial X} - \frac{fV |V| p}{2a}$$

We have eliminated the space and time derivatives of the density ρ in the mass conservation equation using

$$d\rho = \left(\frac{\partial \rho}{\partial H} \right)_P dH + \left(\frac{\partial \rho}{\partial P} \right)_H dP$$

We still must know the density as a function of P and H as well as its derivatives with respect to P and H . The equations are simplified by writing them in terms of the Gruniesen parameter ϕ and the acoustic velocity c :

$$c^2 = 1 / \left(\left(\frac{\partial \rho}{\partial P} \right)_H + \frac{1}{\rho} \left(\frac{\partial \rho}{\partial H} \right)_P \right)$$

$$\phi = -\frac{c^2}{\rho} \left(\frac{\partial \rho}{\partial H} \right)_P$$

The necessary helium properties are parametrized as functions of P and H .⁴

Λ is the heat transfer to the fluid per unit length. These sets of equations are effectively coupled through the heat transfer equation (which is discussed below) and the boundary conditions.

The friction factor is f , and p and a the wetted perimeter and cross-sectional area of the channel. The Cryoloop Simulation parameter list is included as Appendix I.

The model is spatially uniform; the overall cross-sectional area, perimeter, heat inputs, and heat transfer coefficient are independent of x . The active portions of the magnet string have been modeled, i.e., the effective length used here neglects the magnet interconnections.

The boundary conditions selected fix the pressure of both streams and the enthalpy of the single-phase stream at the refrigerator end ($x=0$). The mass flow is fixed at the valve and the mass and enthalpy are continuous across it. The mass flow is chosen so that the fluid is 10% liquid at the exit. The model makes no provision for power or correction lead flows, but includes heat leak due to the power leads.

Initial conditions ($t=0$) must be provided for all values of x . The program is required to converge to a good steady state solution ($\partial(\)/\partial t \rightarrow 0$) as a starting point for the time-dependent studies.

The heat transfer to the fluid, Λ , is composed of two parts: the static load from heat leaks and the ramping load. At each spatial and temporal point, the heat transfer Λ is calculated from the external heat applied to each stream and the heat transfer from the single-phase to the two-phase stream. This stream-to-stream heat transfer is proportional to the temperature difference between the two streams and the heat transfer coefficient, which is a function of the geometry and materials. Although this coefficient does have some flow dependence, the average value used here provides a good approximation for the flow studied.

The equation used for the heat transfer per unit length is

$$Q_T = p_1 \cdot h_{tr} \cdot (T_1 - T_2) / 2$$

where p_1 is the total wetted perimeter of stream 1, and h_{tr} , the heat transfer coefficient, was taken to be $3.35 \text{ mW/cm}^2\text{-}^\circ\text{K}.$ ⁵

In this paper P_1 , H_1 , and V_1 refer to the single-phase stream and P_2 , H_2 , and V_2 refer to the two-phase stream. Similarly, T_1 and T_2 refer to the temperatures, and R_1 and R_2 refer to the densities of the streams.

Steady State

A steady state case has been solved for an average anticipated load: average dynamic load of 250 watts, 10 watts heat leak into the single-phase stream, and 140 watts of heat leak directly into the 2-phase stream.

The boundary conditions at $x=0$ are

$$H1 = 11.43 \text{ j/g}$$

$$P1 = 1.8 \text{ atm}$$

$$P2 = 1.209 \text{ atm}$$

The boundary conditions at $x=L$ are

$$V1 = \dot{M}/R1 \cdot a_1 \text{, where } \dot{M} = 23.81 \text{ g/sec, } R1 \text{ is the density of stream 1 at } x=L, \text{ and } a_1 \text{ is the cross sectional area of stream 1 (19.8 cm}^2\text{).}$$

$$H1 = H2$$

$$\dot{M}_1 = V1/R1 \cdot a_1 = \dot{M}_2 = V2/R2 \cdot a_2 \text{, where } R2 \text{ is the density of stream 2 at } x=L \text{ and } a_2 \text{ is the cross sectional area of stream 2 (11.255 cm}^2\text{).}$$

Fig. 3 shows $T1$, the temperature of stream 1, $T2$, the temperature of stream 2, and Q_T , the heat transferred from stream 1 to stream 2 as a function of the distance along the magnet string.

Notice that $x=0$ is the inlet and $x=L$ the turnaround.

Figs. 4 and 5 show the pressures of the two streams as functions of x . The densities, $R1$ and $R2$, are shown in Figs. 6 and 7. $R1$ shows the simple effects of increasing temperature. $R2$ reflects the increasing proportion of gas from $x=L$ to $x=0$. The velocities of the fluid, $V1$ and $V2$, are shown in Figs. 8 and 9, and the enthalpies, $H1$ and $H2$, in Figs. 10 and 11.

Ramping

The ramping current can be any $I(t)$ for which $Q(t)=F(I(t))$ is known. The ramp model used (200 A/sec, $I_{\max}=4000$ A, 5 sec flat top and 60 sec cycle time) closely parallels anticipated ramping (Fig. 12). The load is taken to be proportional to $(\partial B/\partial t)^n$ and modified to one second rise and fall times for computational convenience. The resultant trapezoidal wave form has a peak of 375 W and an average of 250 W. The heat leak is the same as the steady state case.

The ramp is turned on to the steady state solution. The boundary conditions are the same as the steady state.

The effect of ramping on the temperature is shown in Fig. 13 at $x=0$, L , and $L/2$ for two ramp cycles and both streams. Large values of $\partial(T1)/\partial t$ have correspondingly large $\partial(R1)/\partial t$ that provide velocities of expansion and contraction that subtract from, and add to, the transport velocity. The velocities of the streams at $x=0$, L , and $L/2$ are shown in Figs. 14 and 15. Fixing the flow rate at the valve results in large variations at the refrigerator.

Figs. 16 and 17 show the pressures as a function of t at $x=0$, $L/2$; and L ; Figs. 18 and 19, the enthalpies; and Figs. 20 and 21, the densities.

Ramp Off

The Saver Operations Group empirically demonstrated a cryoloop upset behavior caused by suddenly turning the full ramp off. Under the controls appropriate to the steady state, a dramatic flow imbalance is observed at the refrigerator and it warms up significantly at its cold end. The magnet simulator was tested, in a preliminary way, to study that situation. The 250 watt steady state dynamic load was reduced to 10 watts in one second, and the boundary conditions held to approximate the "frozen" controls on ramp-off that are now in effect.

The temperature distributions as a function of x at 20 second intervals are shown in Figs. 22 and 23. The $t=0$ curves correspond to the steady state temperature distribution (see Fig. 3). Figs. 24 through 32 show the behavior of the other quantities as a function of x at 20 second intervals.

The velocities of the single-phase stream at $x=0$, $L/2$, and L are shown in Fig. 33 as a function of time for 1200 seconds. The velocities rise instantaneously, greater the nearer the inlet, from values that were nearly equal (to within a few percent) at $t=0$. The effect is brought about by the rapid rise in single-phase fluid density in the absence of a heat load. As the stream-to-stream ΔT decays, the effect subsides.

The behavior of the single-phase stream as a function of time, odd-numbered Figs. 33 to 41, describe the situation very well at small t . The actual behavior deviates from

that simulated as t grows larger because of the assumption of constant mass flow. The demand for liquid of the single-phase stream (up to 300 l/hr) at ramp-off cannot be provided by the current refrigerator. The instantaneous peak for the whole Saver would be $48 \times (300 \text{ l/hr}) = 14,400 \text{ l/hr}$.

The behavior of the two-phase stream, even-numbered Figs. 34 through 42, is a good description through the first tens of seconds. The enthalpy, H_2 , Fig. 40, makes it clear that the liquid fraction is increasing dramatically, especially at the exit ($x=0$), after the sharp reduction in load. The real situation differs in two respects at large t : first, the assumption of constant mass flow fails as before, and, second, the integrated effect, the collection of liquid, is enhanced by continued inactivity of the valve.

This indication of the magnitude of the transient behavior has suggested hardware modifications aimed at significantly reducing the effects. These considerations are properly the subject of a paper themselves and are mentioned here in support of the continued development of this work. Ramp-off will, unavoidably, be frequently exercised in the early trials and the cryolooop should be prepared to deal with its consequences. It should be noted that the transition from ramping to extended flat-top, beam coasting, or beam storage mode are different from the ramp-off transient only in the instantaneous ramp rate and the recent ramp history. Item 4 of the work agenda address this concern.

Future Work

Although the work is far from complete, we felt an interim report should be distributed to establish capability, direction, and to solicit comments.

The agenda of work currently is to

- 1) develop a no-ramp-load steady state,
- 2) apply the ramp to this steady state (i.e., saver turn-on) at various rates,
- 3) turn the ramp off,
- 4) from the steady state in item 1, ramp up, then simulate beam storage with flat-tops of various lengths,

5) install two-phase superheat feedback to a mass flow control valve and test the ramp-off behavior.

References

- 1 Jones, M.C., "Cool-down of superconducting power transmission lines with single phase helium," Cryogenics, March 1980, p 139.
- 2 Madsen, N.K., Sincovec, R.F. in Numerical Methods for Differential Systems, L. Lapidus, W. E. Shiesser, ed. (Academic Press Inc., New York, 1976).
- 3 Arp, V., "Thermodynamics of single-phase one-dimensional fluid flow," Cryogenics, May 1975, p 285.
- 4 Arp, V., N.B.S., Boulder, Colo., private communication.
- 5 Mulholland, George, Magnet Heat Transfer, Fermilab Memo, June 25, 1980.

APPENDIX I

APRIL 1981

PARAMETER LIST: CRYOLOOP SIMULATION		VALUE
CHL	1. CHL FEED ENTHALPY (J/g), $N = 0$	12.33
	2. CHL FEED ENTHALPY (J/g), $F(N)$ (SEE NOTE 1)	$P1+\Delta H_2$
	3. CHL PRODUCTION RATE (g/s).	139
	4. FRACTION OF CHL PRODUCTION UTILIZED, X	0.9
LN ₂	5. LN ₂ FEED ENTHALPY (J/g), $N = 0$	34.8
	6. LN ₂ FEED ENTHALPY (J/g), $F(N)$ (SEE NOTE 2)	$P5+\Delta H_5$
	7. LN ₂ FEED PRODUCTION RATE (g/s)	568
	8. FRACTION OF LN ₂ PRODUCTION UTILIZED, Y	0.9
LINES	9. TOTAL HEAT LEAK LHE TRANSFER LINE (J/s).	300
	10. TOTAL HEAT LEAK LN ₂ TRANSFER LINE (J/s).	4500
COMP.	11. TOTAL RING COMPRESSOR FLOW (g/s)/24 (NOMINAL) (SEE NOTE 3).	50
	12. SUCTION MANIFOLD PRESSURE (ATM).	1.03
	13. DISCHARGE MANIFOLD PRESSURE (ATM).	20.0
	14. DISCHARGE TEMPERATURE (°K)	316
LEADS	15. DIPOLE LEAD FLOW TOTAL (g/s)	0.48
	16. CORRECTION LEAD FLOW TOTAL (g/s)	0.48
	17. DIPOLE LEAD MAXIMUM RAMP HEAT LEAK TOTAL	10
	18. CORRECTION LEAD HEAT LEAK TOTAL	10
REFRIG. & ENGS.	19. CRYOLOOP REFRIGERATOR NUMBER, N	24
	20. STATIC HEAT LEAK (at 5-6°K), REFRIG. (J/s)	10
	21. STATIC HEAT LEAK, GAS ENGINE (J/s)	10
	22. GAS ENGINE DISPLACEMENT (cm ³).	463
	23. MAXIMUM GAS ENGINE SPEED (RPM)	317
	24. MINIMUM GAS ENGINE SPEED (RPM)	120

CONT'D.

<u>PARAMETER LIST: CRYOLOOP SIMULATION (CONT'D.)</u>		<u>VALUE</u>
	25. EFFICIENCY GAS ENGINE	0.65
	26. INTAKE CUT-OFF ENGINE	0.38
	27. STATIC HEAT LEAK WET ENGINE (J/s)	16
	28. WET ENGINE DISPLACEMENT (cm ³)	206
	29. WET ENGINE MAXIMUM SPEED (RPM).	317
	30. WET ENGINE MINIMUM SPEED (RPM).	107
	31. EFFICIENCY OF WET ENGINE (NOMINAL).	0.7
	32. INTAKE CUT-OFF WET ENGINE	0.7
	33. SHELL SIDE PRESSURE DROP (ATM).	0.195
	34. TUBE SIDE PRESSURE DROP (ATM)	1.0 TYP
	35. SINGLE PHASE SOURCE PRESSURE (ATM).	1.8
	36. 20 RETURN PRESSURE (ATM), P12+P33	1.23 TYP
	37. HEAT EXCHANGER CAPACITY (W/°K), (W/40) ^{°-°} (1.4×10 ⁴)	1.67E4
SUB-	38. SINGLE PHASE SOURCE TEMPERATURE (°K).	4.5 TYP
COOLER	39. SUPER HEAT OF 20 RETURN (°K)	0.1 TYP
	40. SUBCOOLER TEMPERATURE APPROACH (°K)	0.09
	41. SUBCOOLER RESERVOIR VOLUME (cm ³).	10E3
MAG.	42. NUMBER OF DIPOLES IN CRYOLOOP	16
HALF	43. NUMBER OF QUADRUPOLES IN CRYOLOOP	4
STRING	44. STATIC HEAT LEAK/DIPOLE (J/s)	7
	45. STATIC HEAT LEAK/QUADRUPOLE (J/s)	7
	46. HALF TOTAL OF OTHER STRING HEAT LEAKS (J/s)	20
	47. 10 DIPOLE CROSS SECTIONAL AREA (cm ²), 7.58+12.22	19.8
	48. 20 DIPOLE CROSS SECTIONAL AREA (cm ²).	11.3
	49. 10 QUADRUPOLE CROSS SECTIONAL AREA (cm ²), 17.3+7.20.	24.5
	50. 20 QUADRUPOLE CROSS SECTIONAL AREA (cm ²).	16.5
	51. 10 DIPOLE WETTED PERIMETER (cm)	119.3

CONT'D.

<u>PARAMETER LIST: CRYOLOOP SIMULATION (CONT'D.)</u>		<u>VALUE</u>
52. 20 DIPOLE WETTED PERIMETER (CM)		87.9
53. 10 QUADRUPOLE WETTED PERIMETER (CM), 50.2+95.7.		146
54. 20 QUADRUPOLE WETTED PERIMETER (CM), (ESTIMATE)		90
55. DIPOLE LENGTH (CM)		608
56. QUADRUPOLE LENGTH (CM)		175
57. SUBCOOLER-LHE ENGINE POTENTIAL (CM).		914
58. RAMPING LOAD, AVERAGE (J/s)		240
59. FRICTION FACTOR 10		5.26E-3
60. FRICTION FACTOR 20 AVERAGE		4.45E-3
61. LOCAL ATMOSPHERIC PRESSURE (A _S).		0.973
62. HEAT TRANSFER COEF. (WINDING TO 20) MW/CM°K.		3.35 TYP

NOTES:

$$1. \quad \Delta H_2 = \sum_{i=0}^{i=N} \frac{P_9}{(24-i(P_4))(P_3)}$$

$$2. \quad \Delta H_5 = \sum_{i=0}^{i=N} \frac{P_{10}}{(24-i(P_8))(P_7)}$$

3. MASS DEPENDENT PARAMETER VALUES ARE PROVIDED AT THE FLOW RATE OF PARAMETER 11.

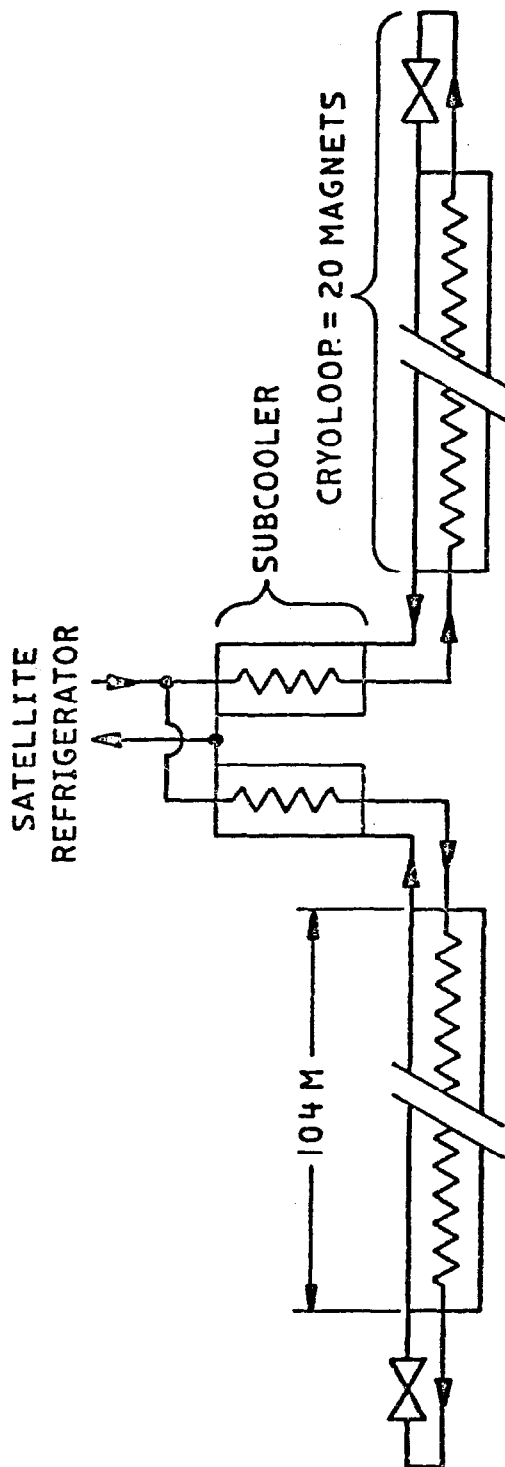


Figure 1. A Satellite drives parallel cryoloops.

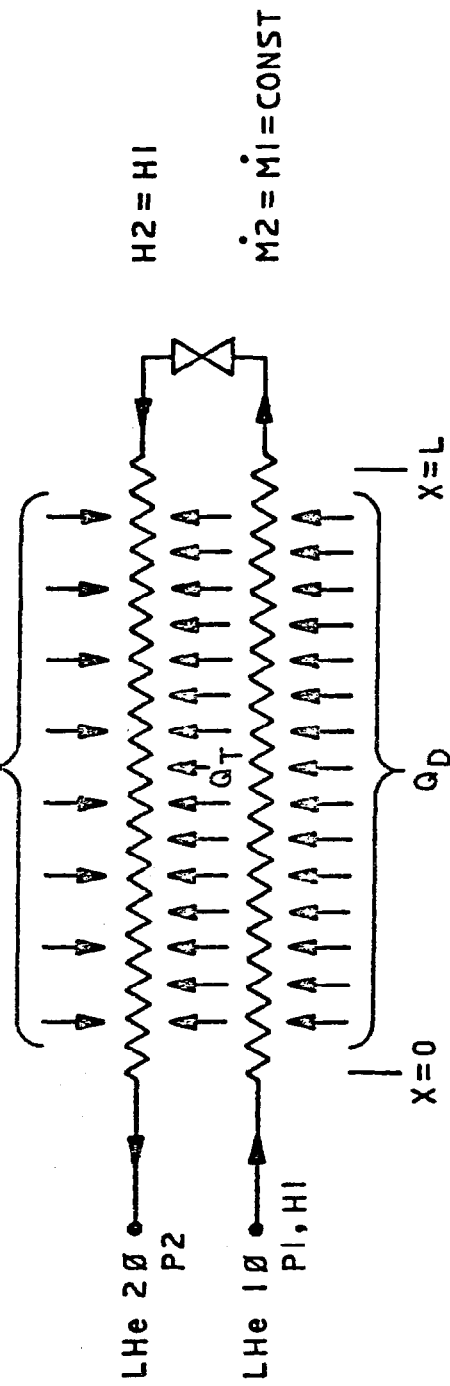


Figure 2. A cryoloop as a distributed heat exchanger.

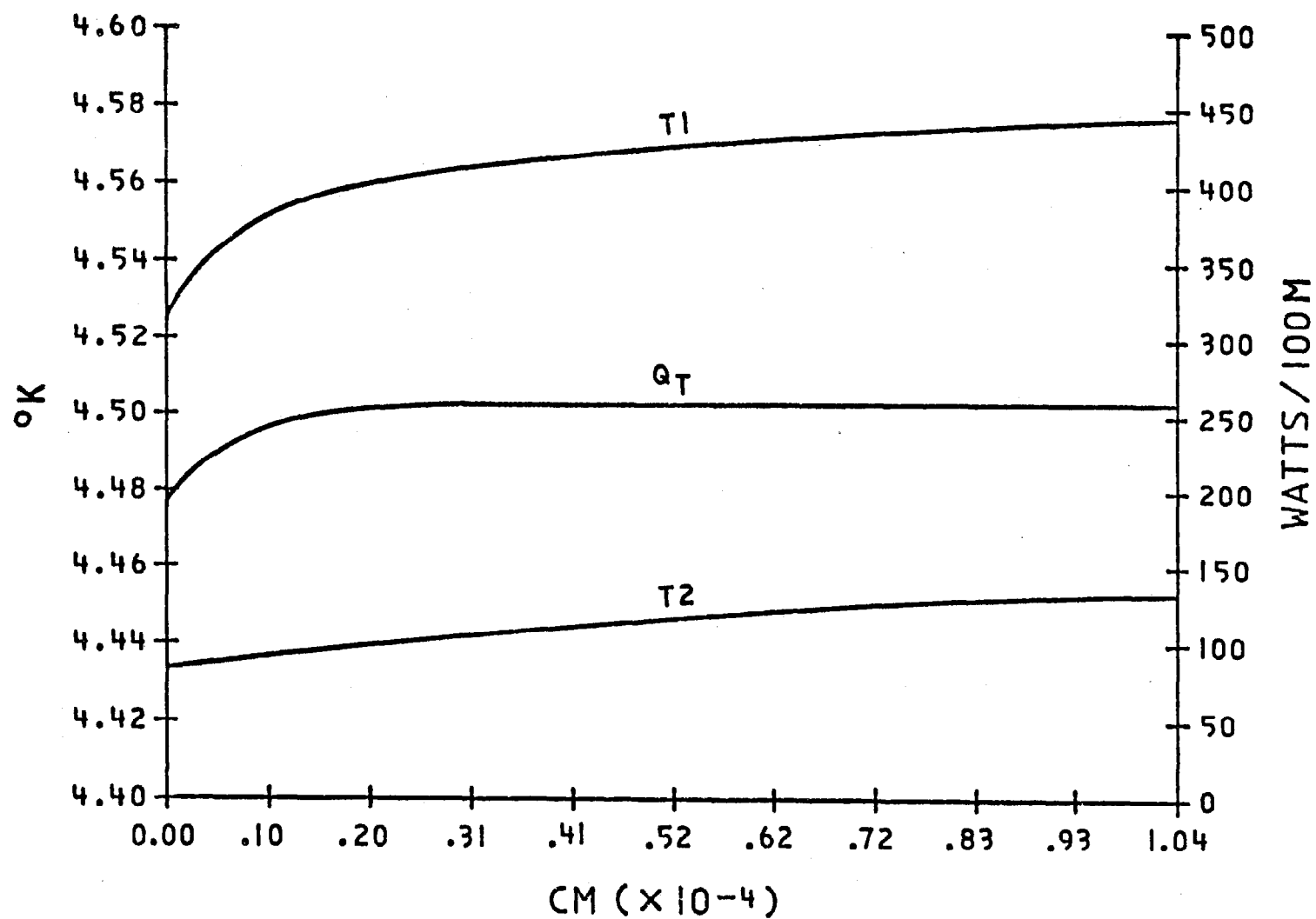


Figure 3. Temperatures and heat transfer of steady state (DC) mode.

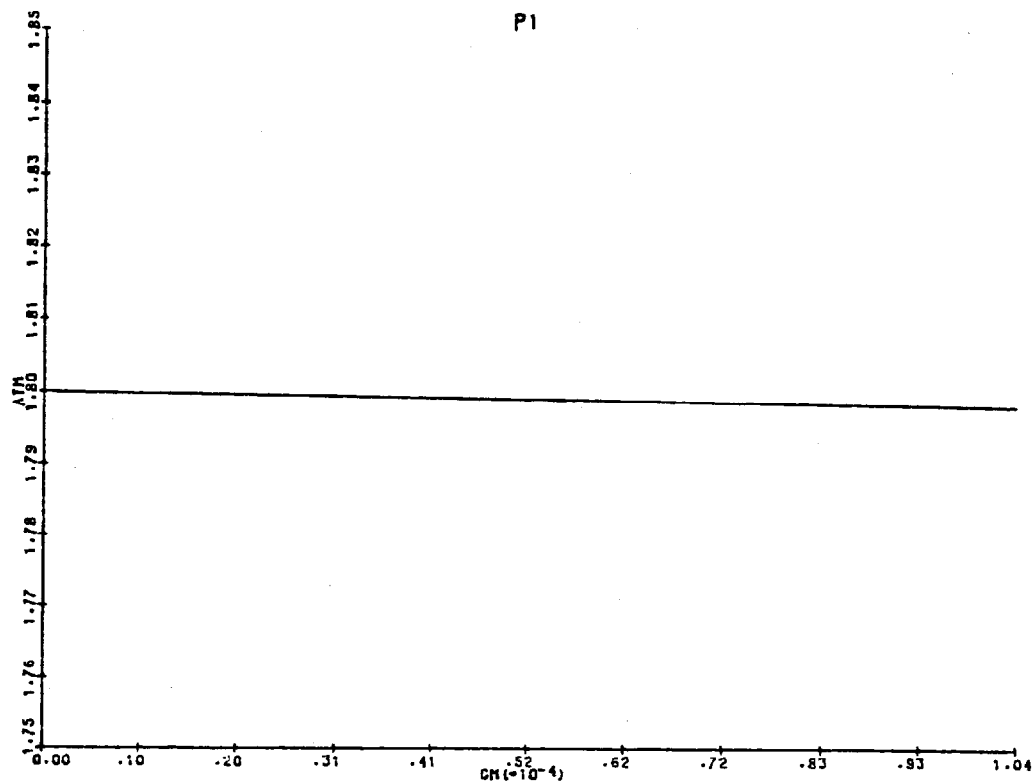


Figure 4. Pressure of single-phase stream as a function of x .

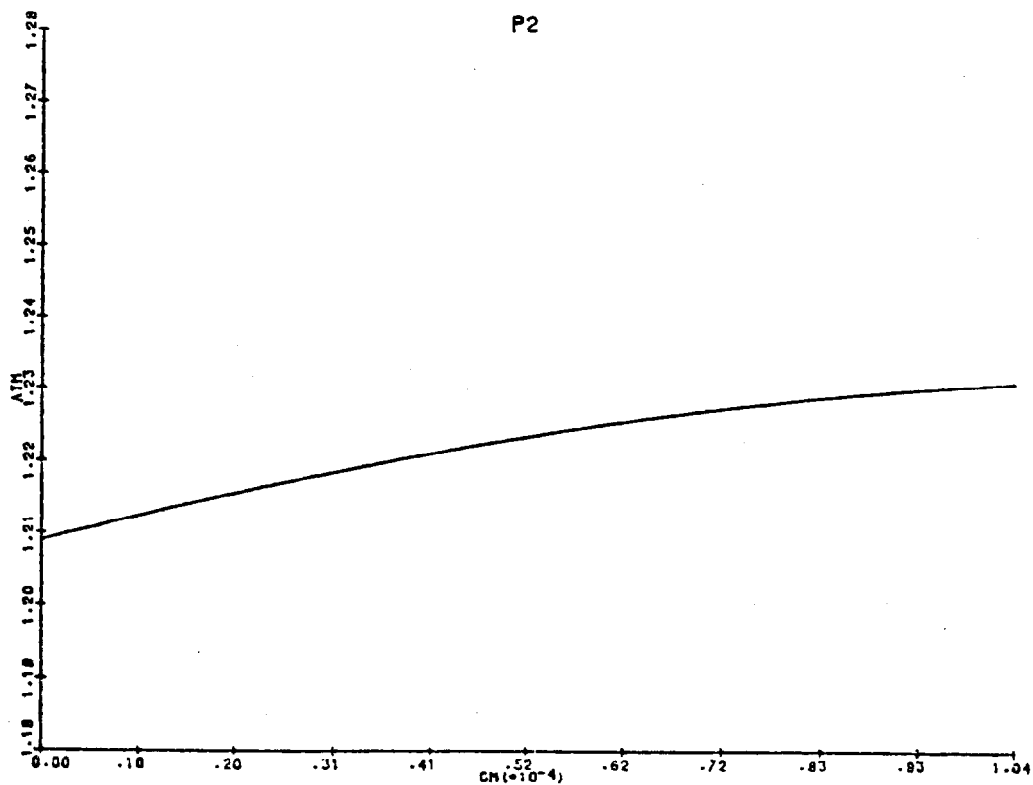


Figure 5. Pressure of the two-phase stream as a function of x .

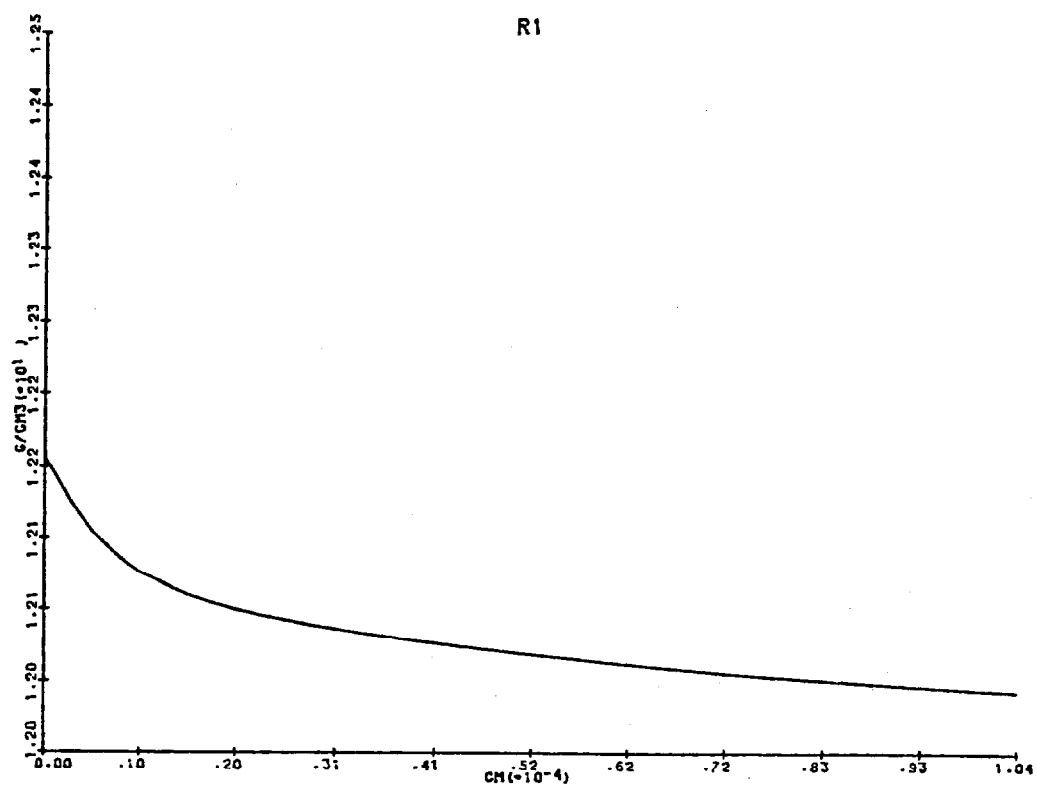


Figure 6. Density of single-phase stream as a function of x .

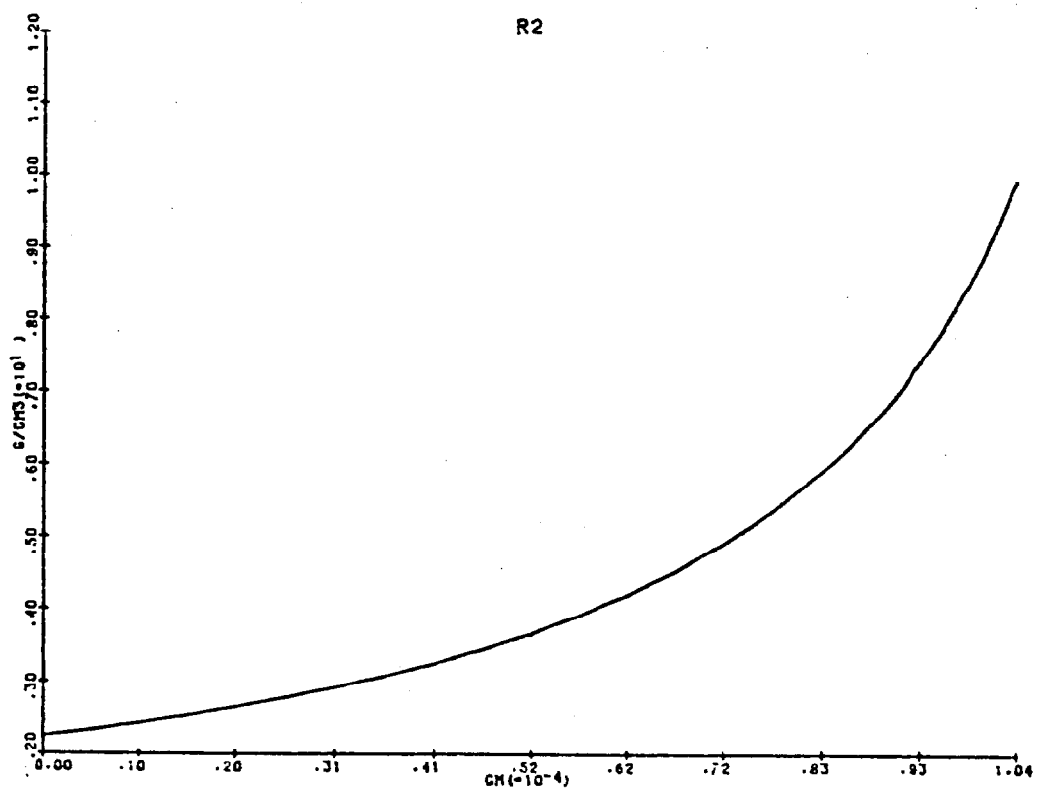


Figure 7. Density of two-phase stream as a function of x .

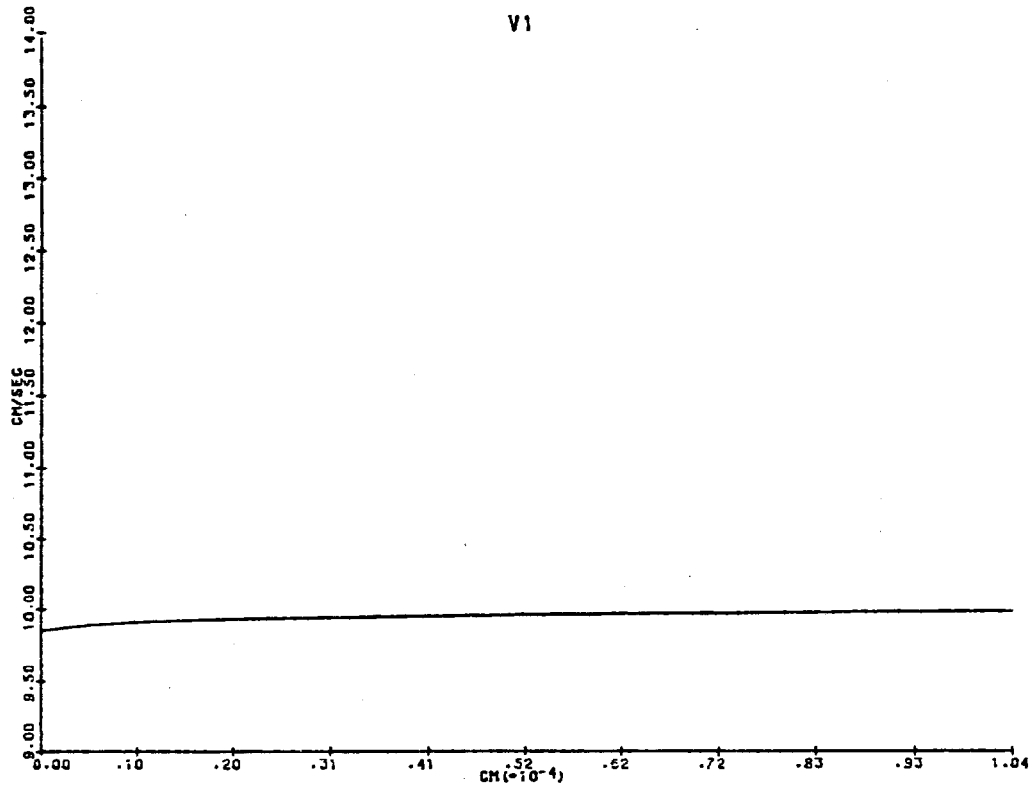


Figure 8. Velocity of single-phase stream as a function of x .

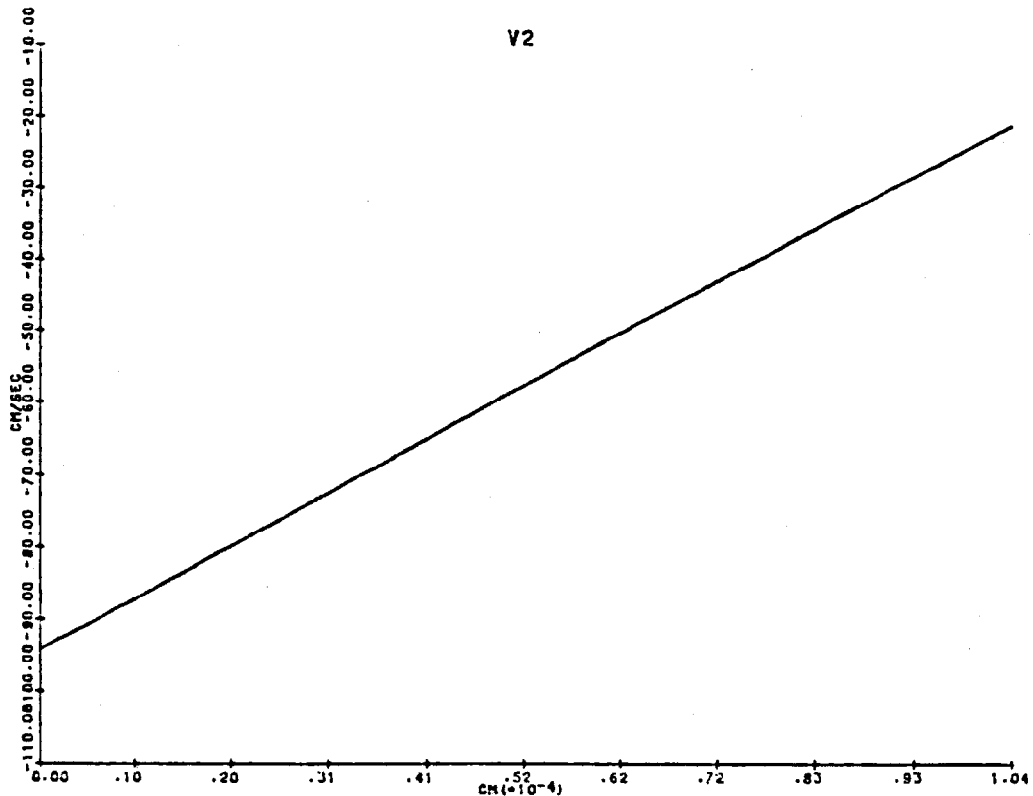


Figure 9. Velocity of two-phase stream as a function of x .

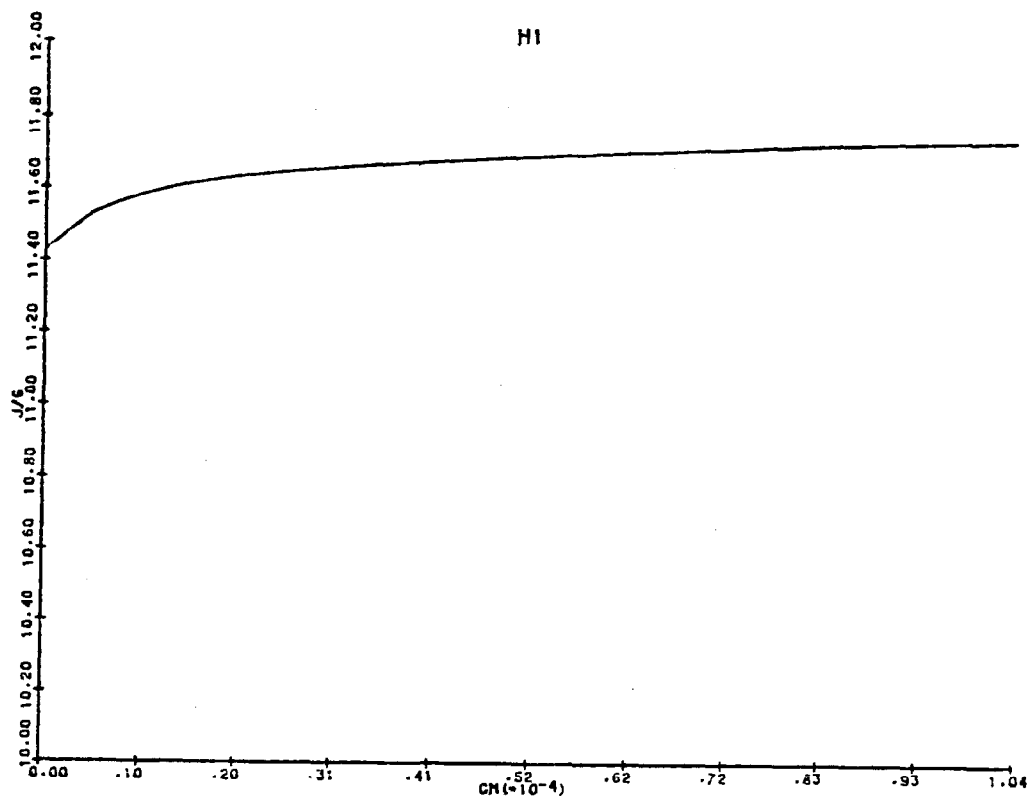


Figure 10. Enthalpy of single-phase stream as a function of x .

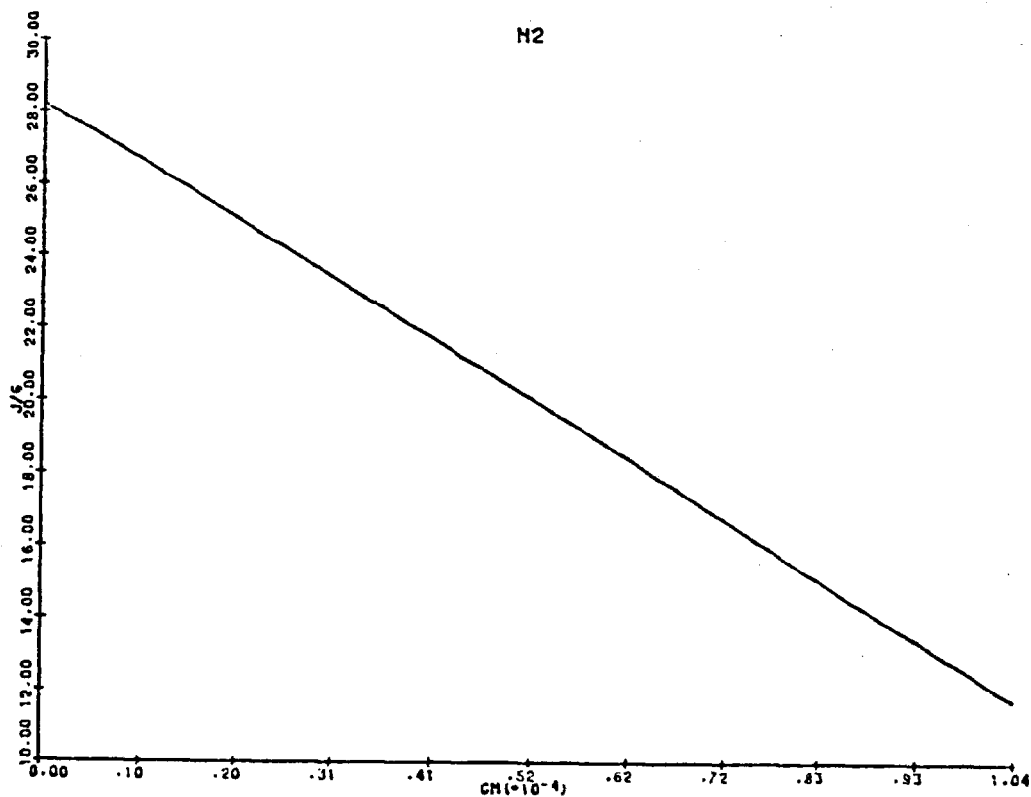
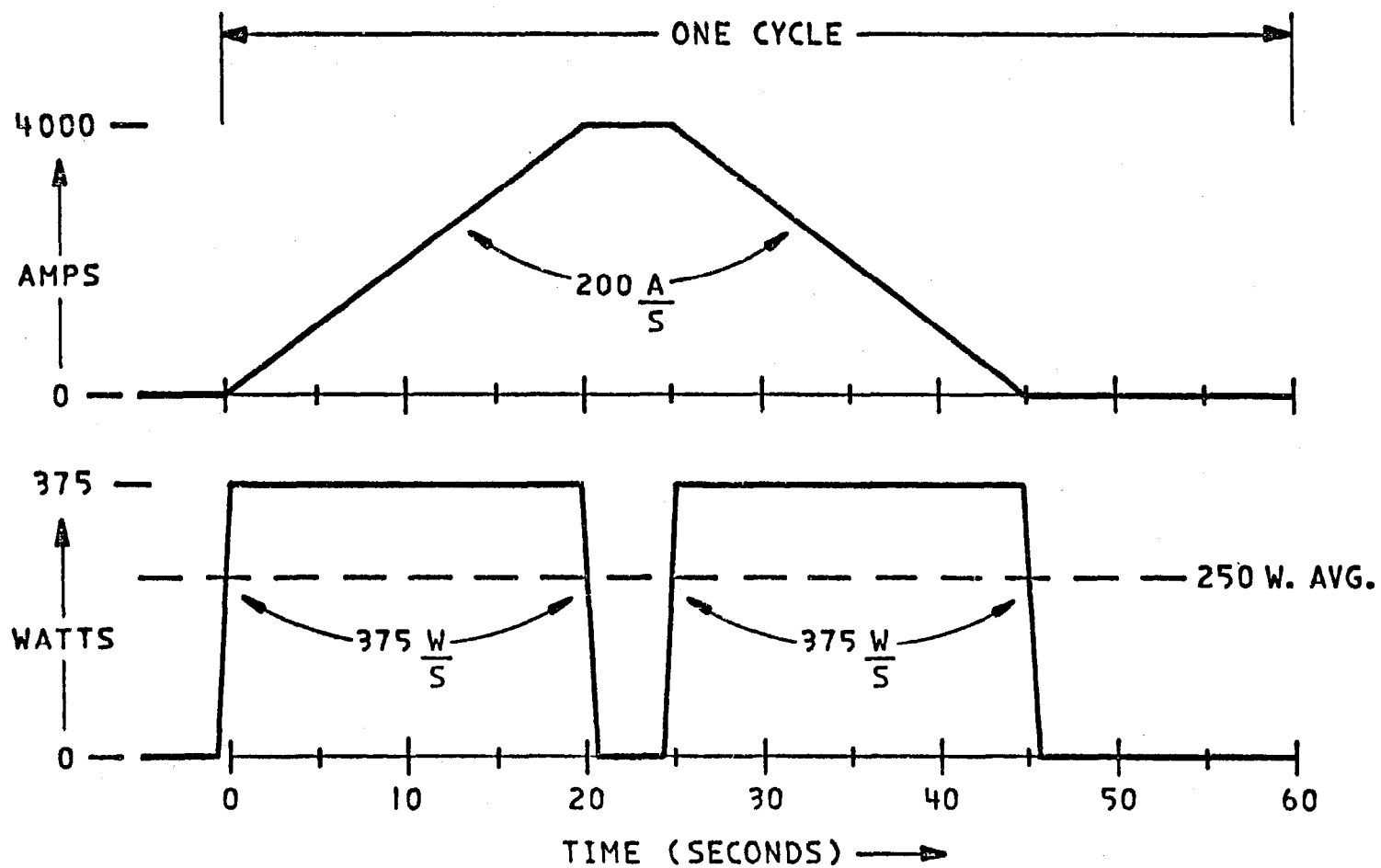


Figure 11. Enthalpy of two-phase stream as a function of x .



MAGNET CURRENT, RAMP LOAD = $F(t)$

Figure 12. Assumed ramping current and refrigeration load as a function of time.

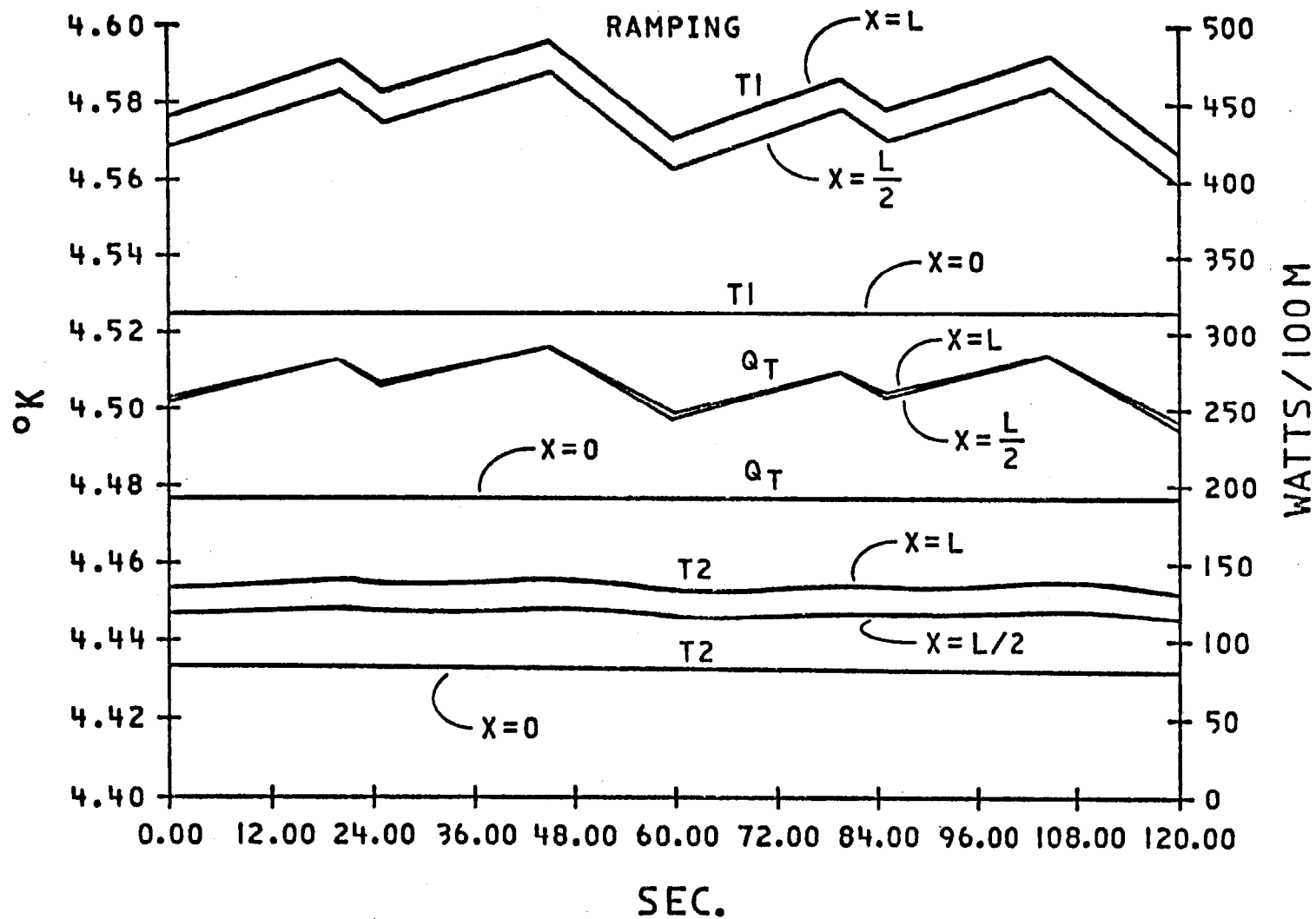


Figure 13; Heat transfer and single- and two-phase temperatures for two ramp cycles.

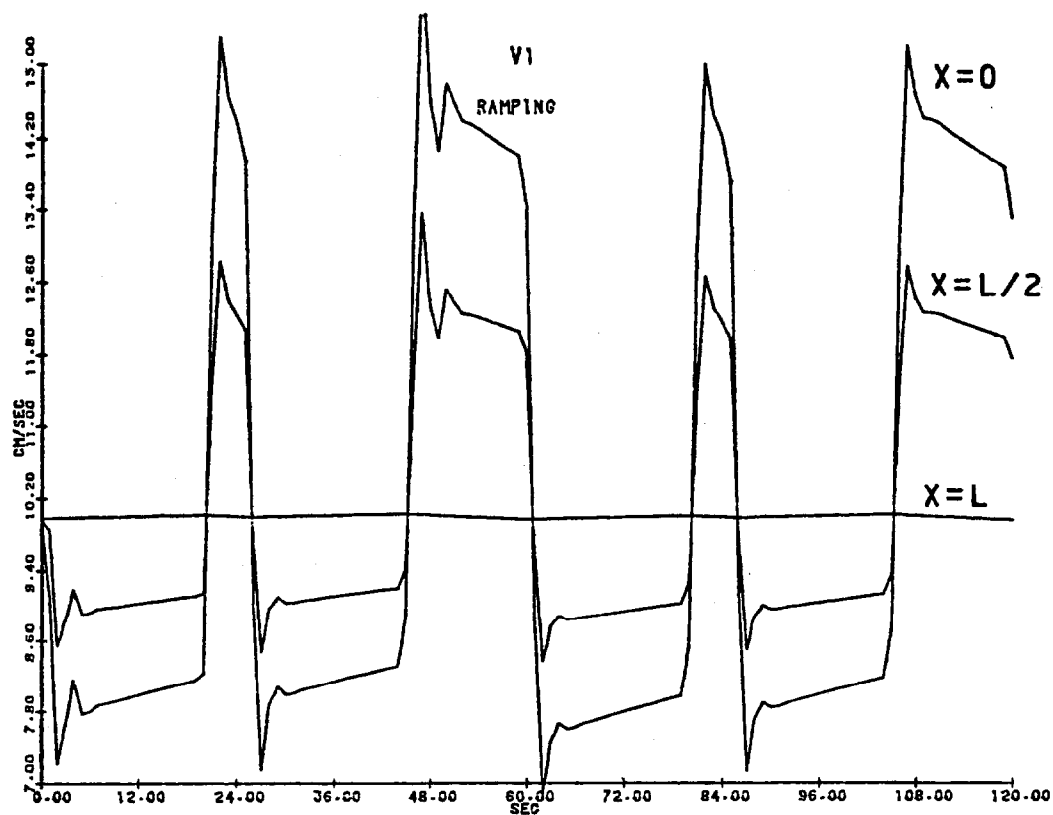


Figure 14. Velocity of single-phase stream over two ramp cycles.

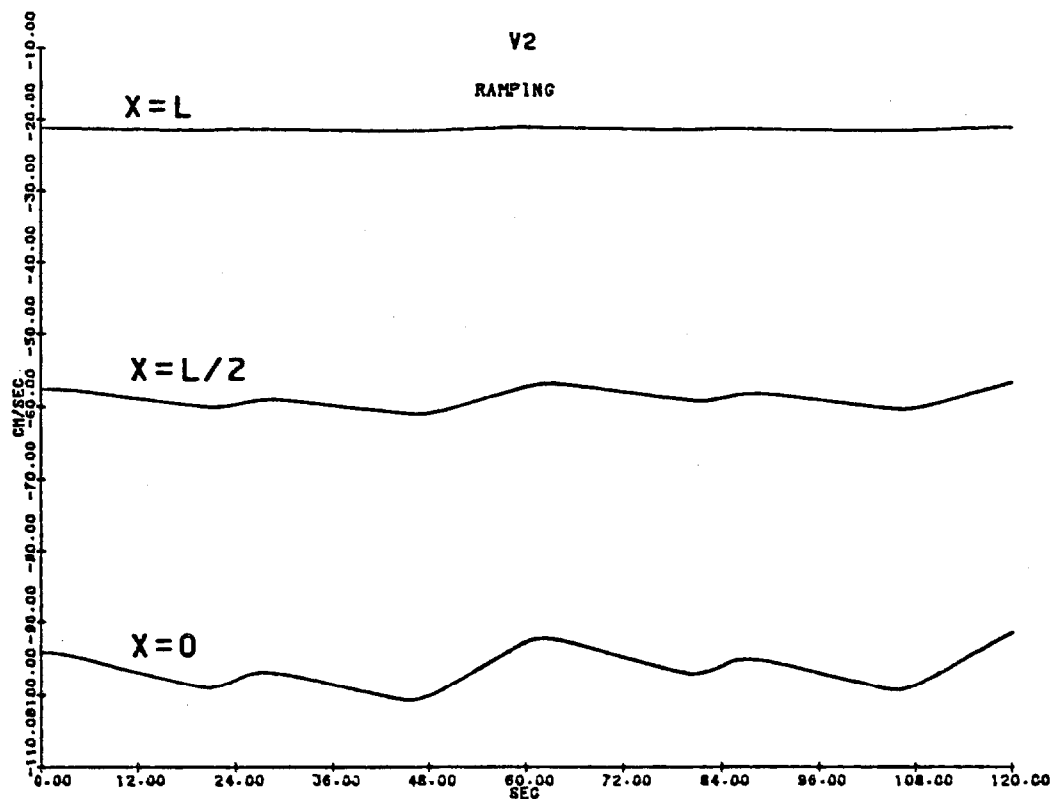


Figure 15. Velocity of two-phase stream over two ramp cycles.

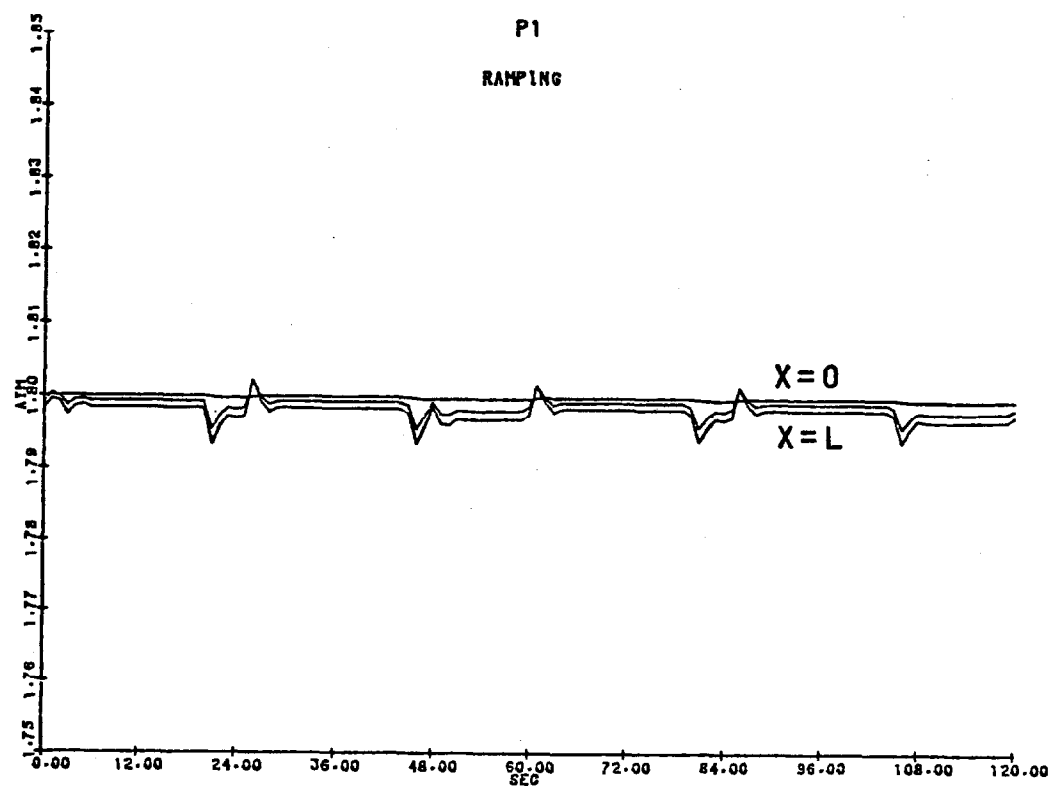


Figure 16. Pressure of single-phase stream over two ramp cycles.

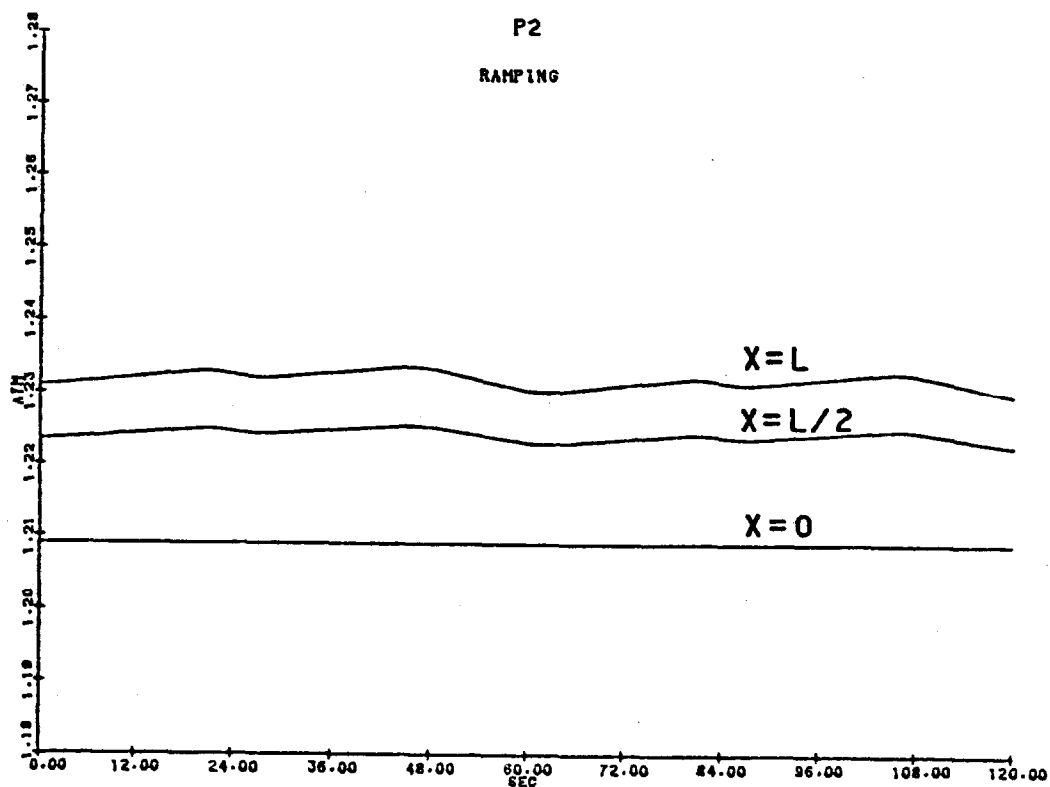


Figure 17. Pressure of two-phase stream over two ramp cycles.

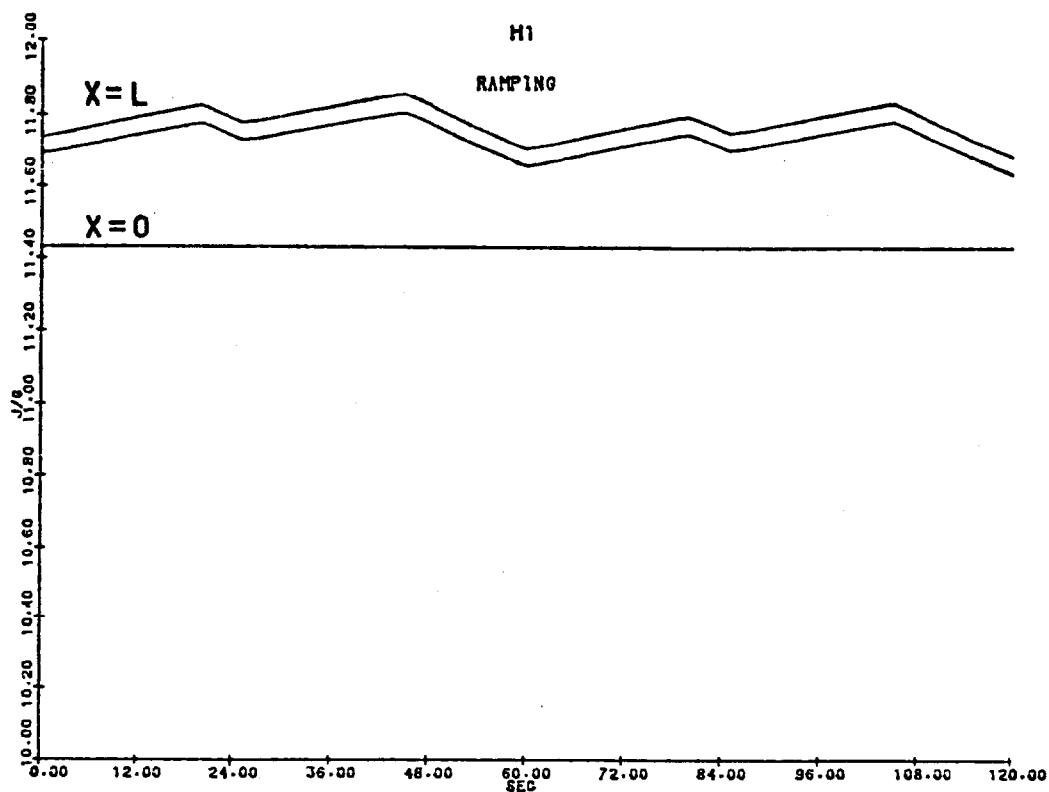


Figure 18. Enthalpy of single-phase stream over two ramp cycles.

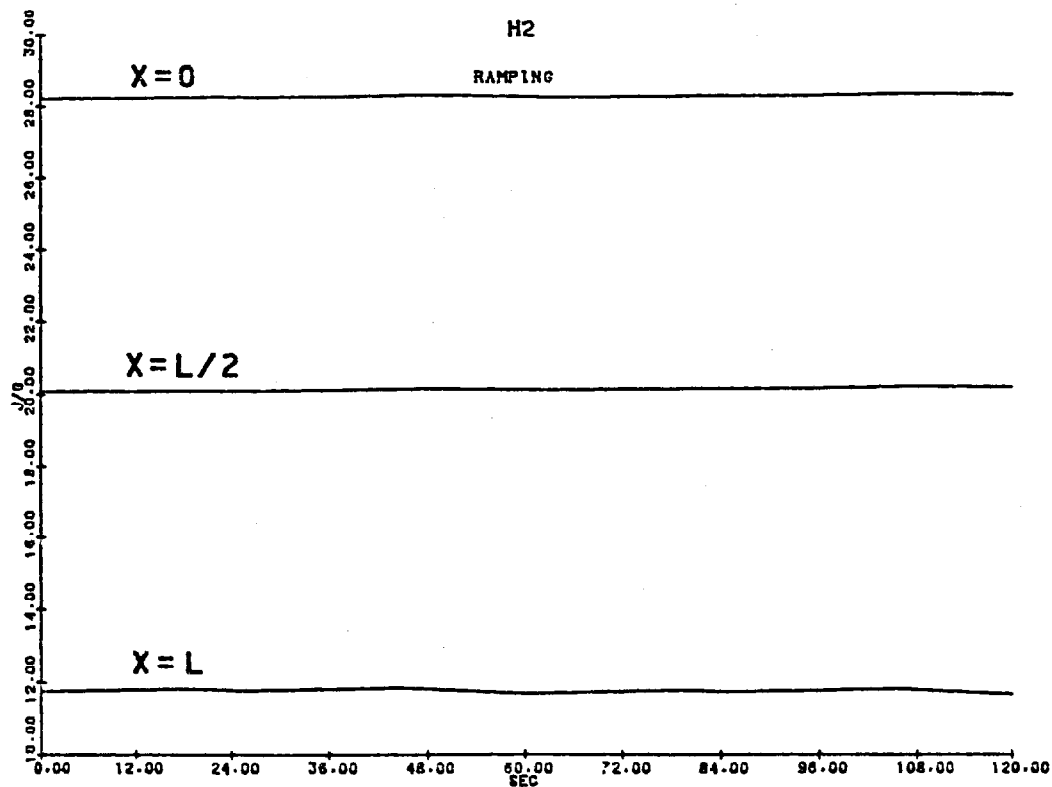


Figure 19. Enthalpy of two-phase stream over two ramp cycles.

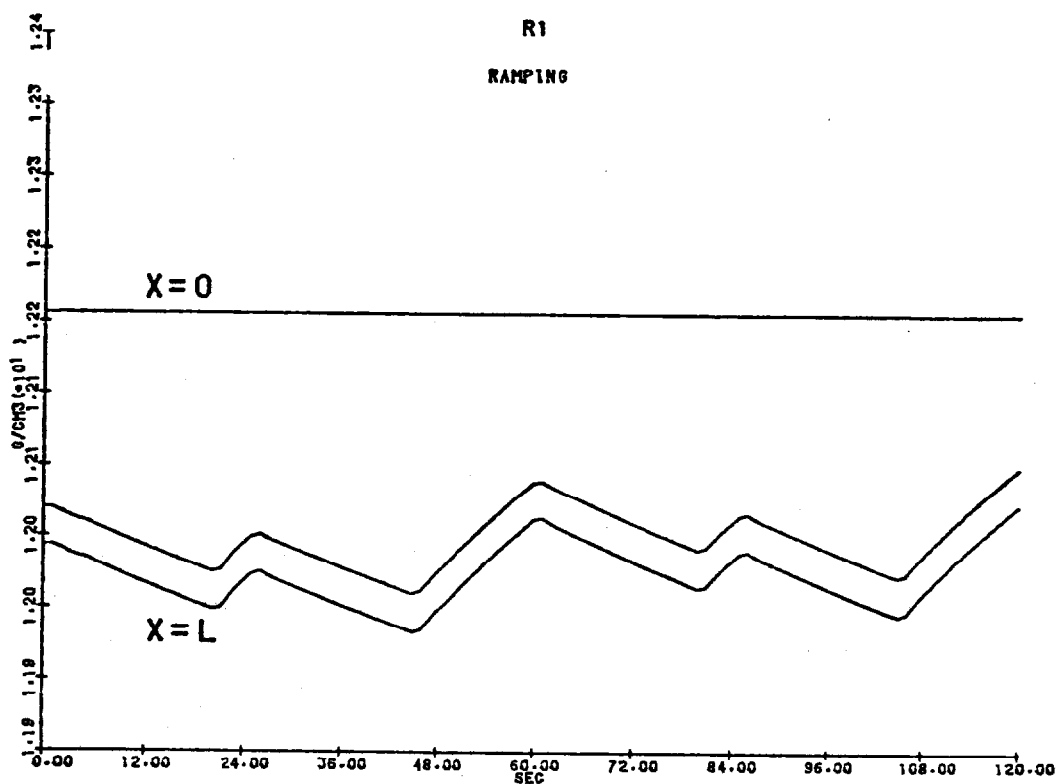


Figure 20. Density of single-phase stream over two ramp cycles.

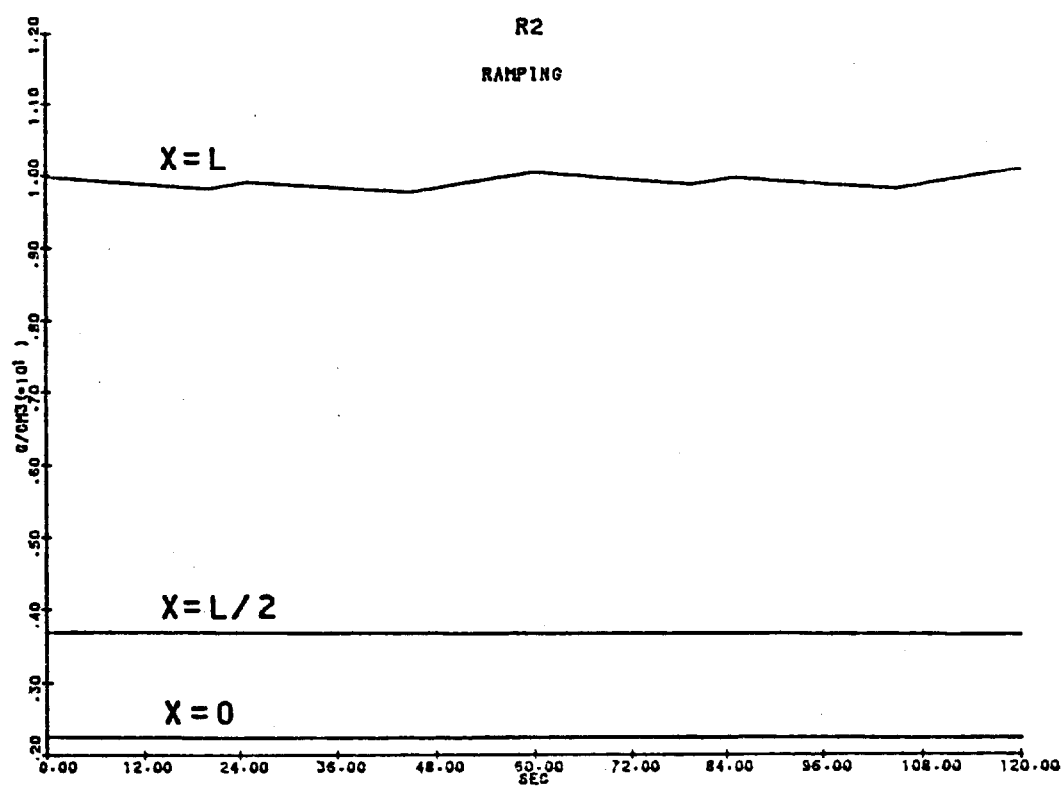


Figure 21. Density of two-phase stream over two ramp cycles.

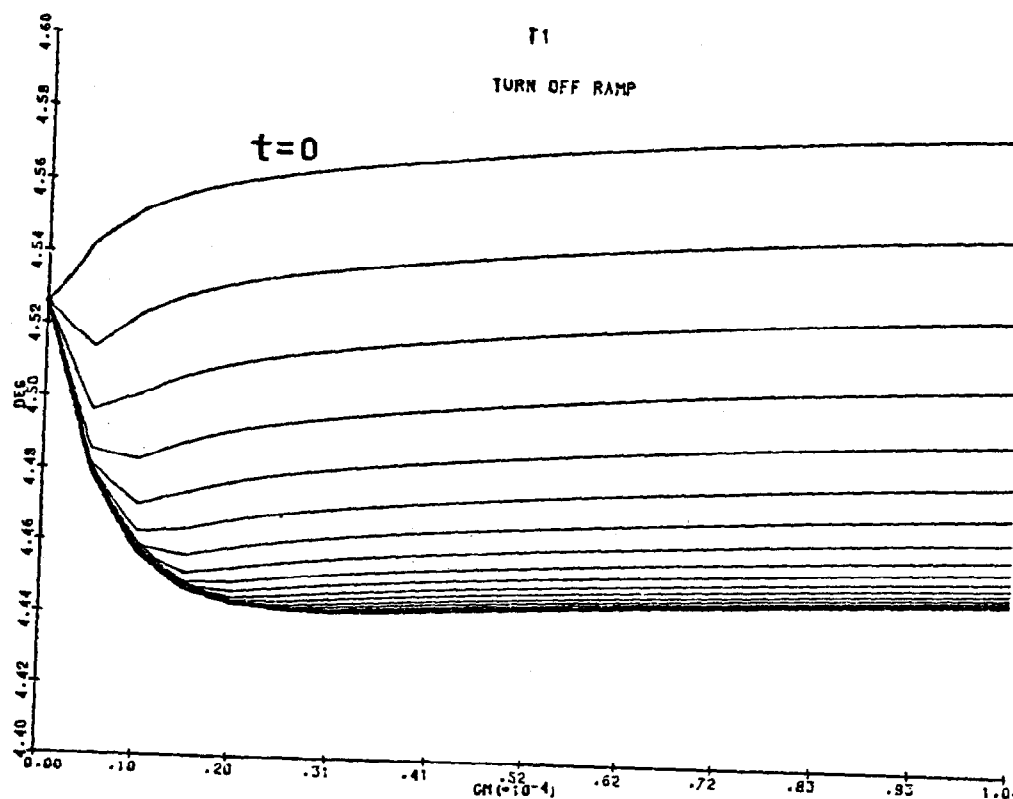


Figure 22. Temperature distribution of single-phase stream as a function of x at 20 sec intervals.

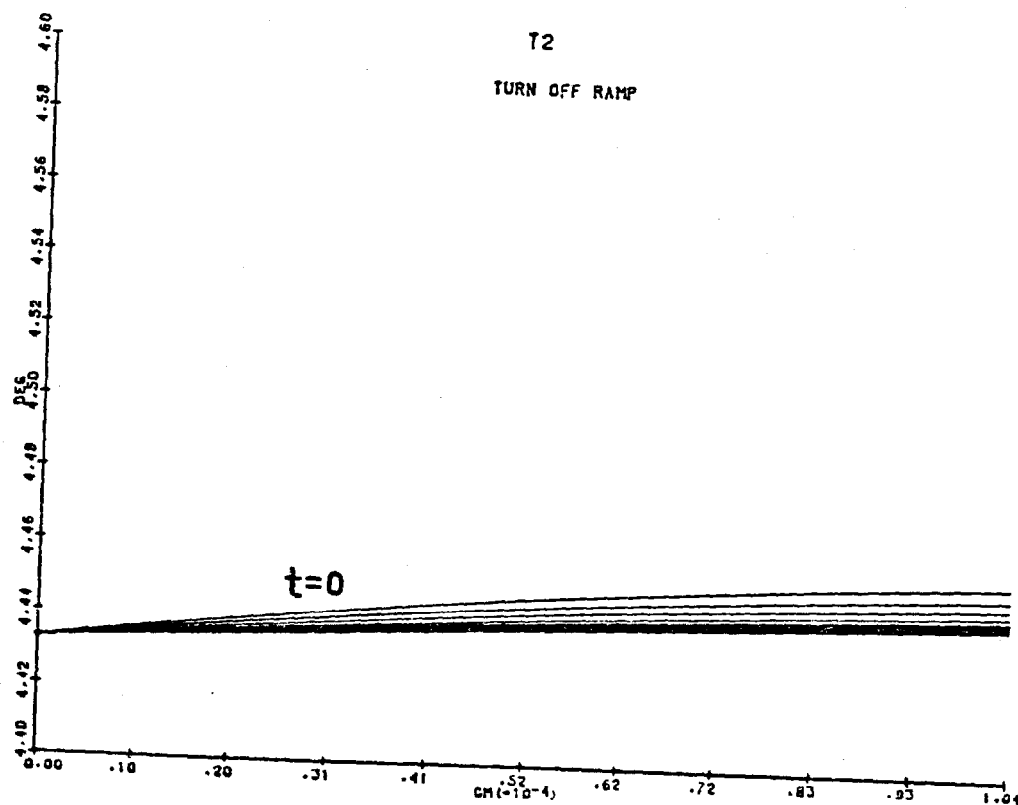


Figure 23. Temperature distribution of two-phase stream as a function of x at 20 sec intervals.

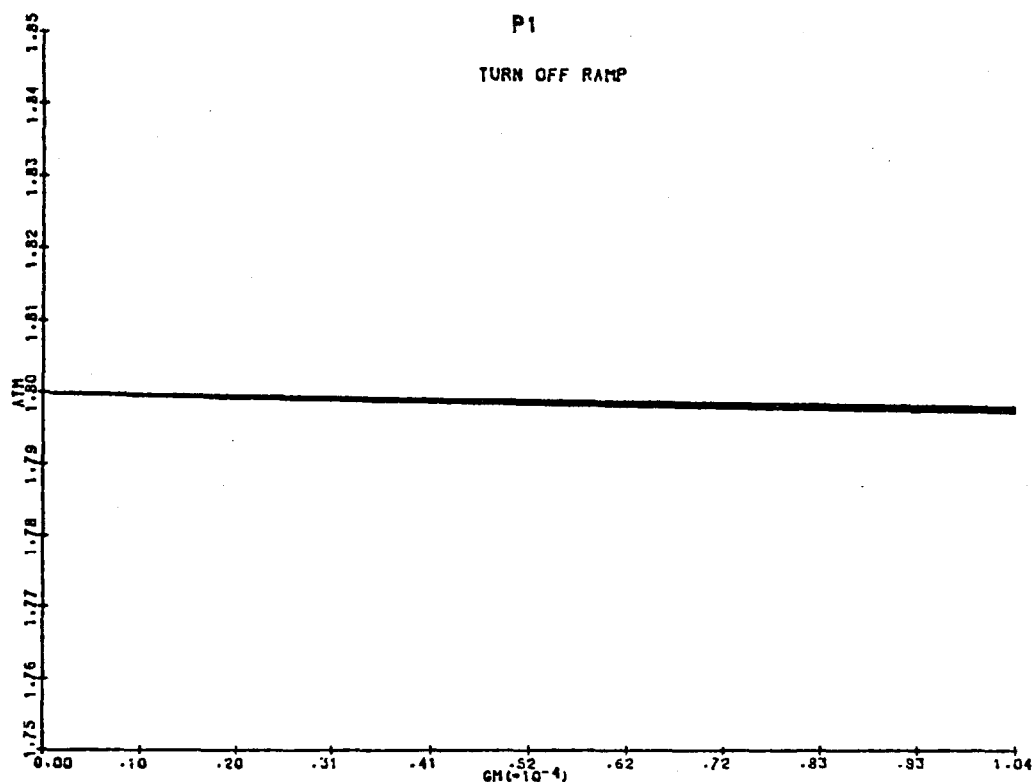


Figure 24. Pressure of single-phase stream as a function of x at 20 sec intervals.

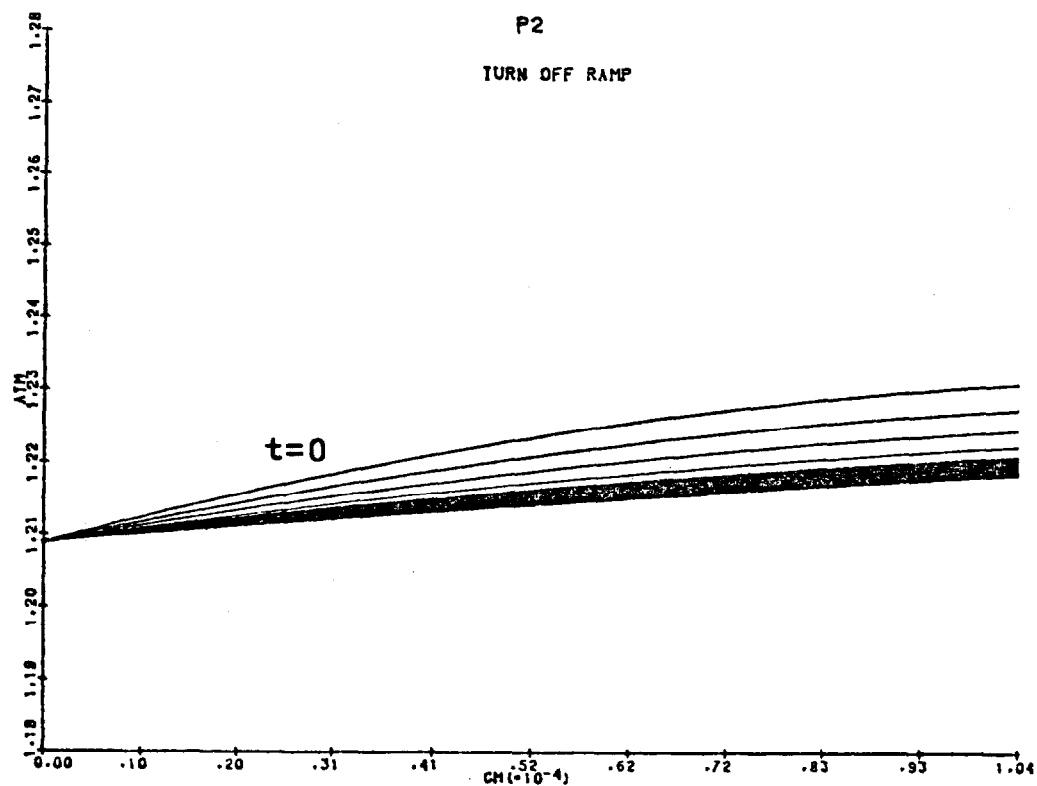


Figure 25. Pressure of two-phase stream as a function of x at 20 sec intervals.

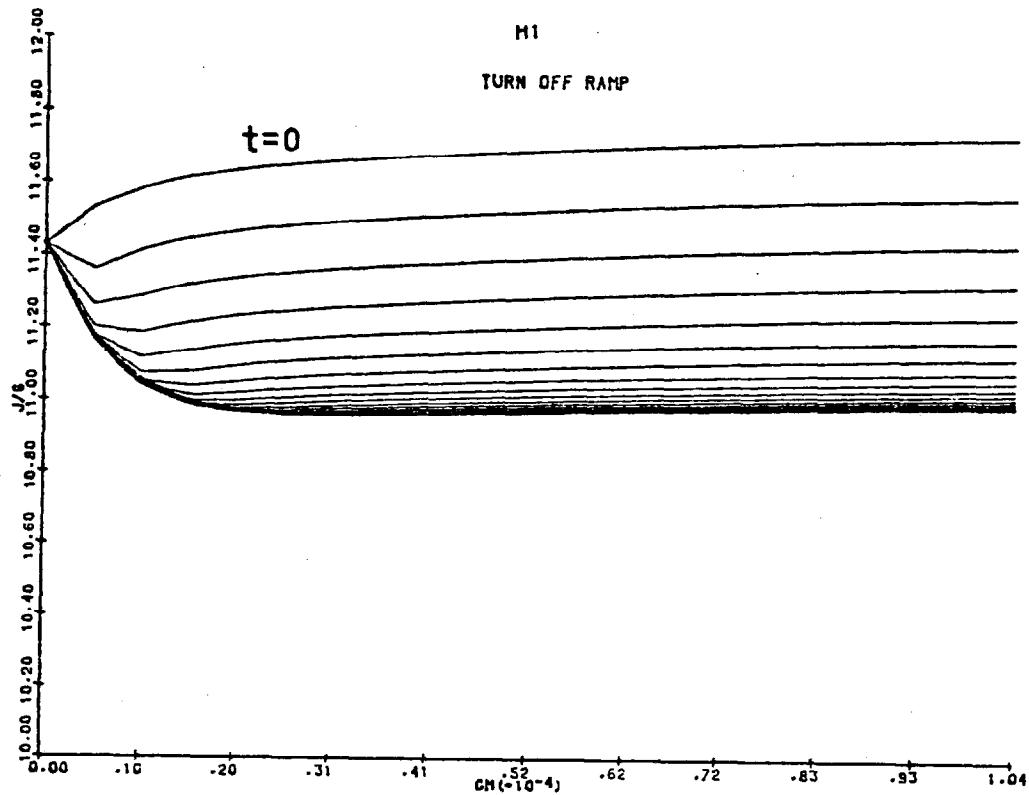


Figure 26. Enthalpy of single-phase stream as a function of x at 20 sec intervals.

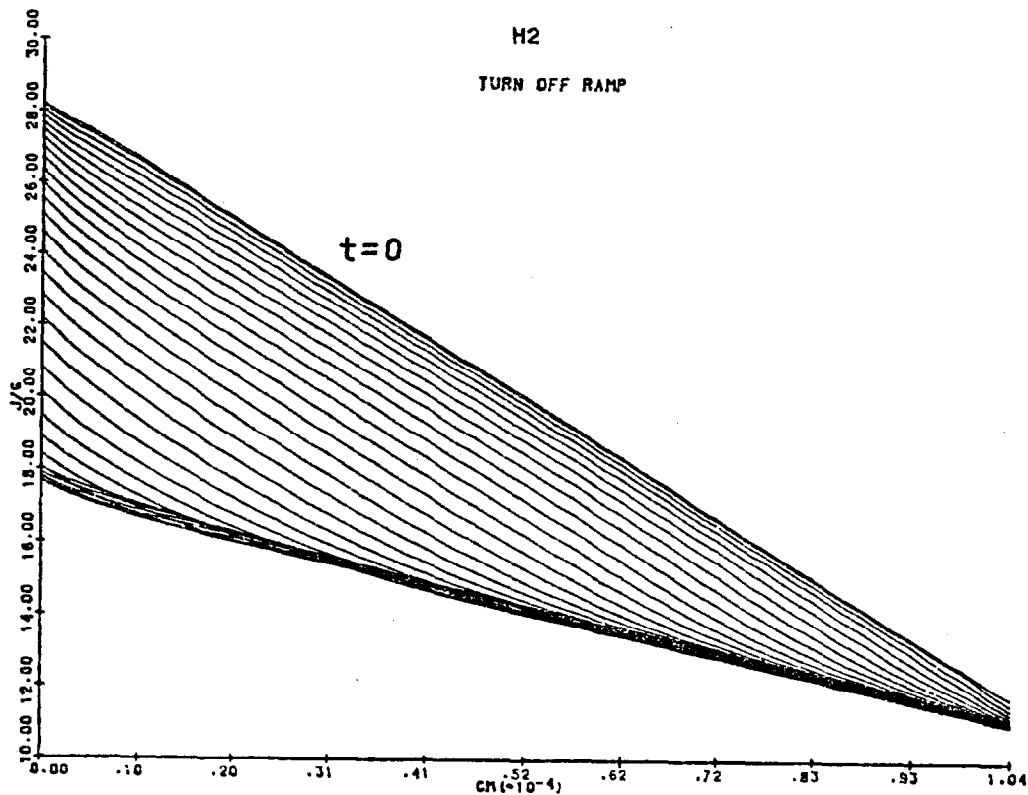


Figure 27. Enthalpy of two-phase stream as a function of x at 20 sec intervals.

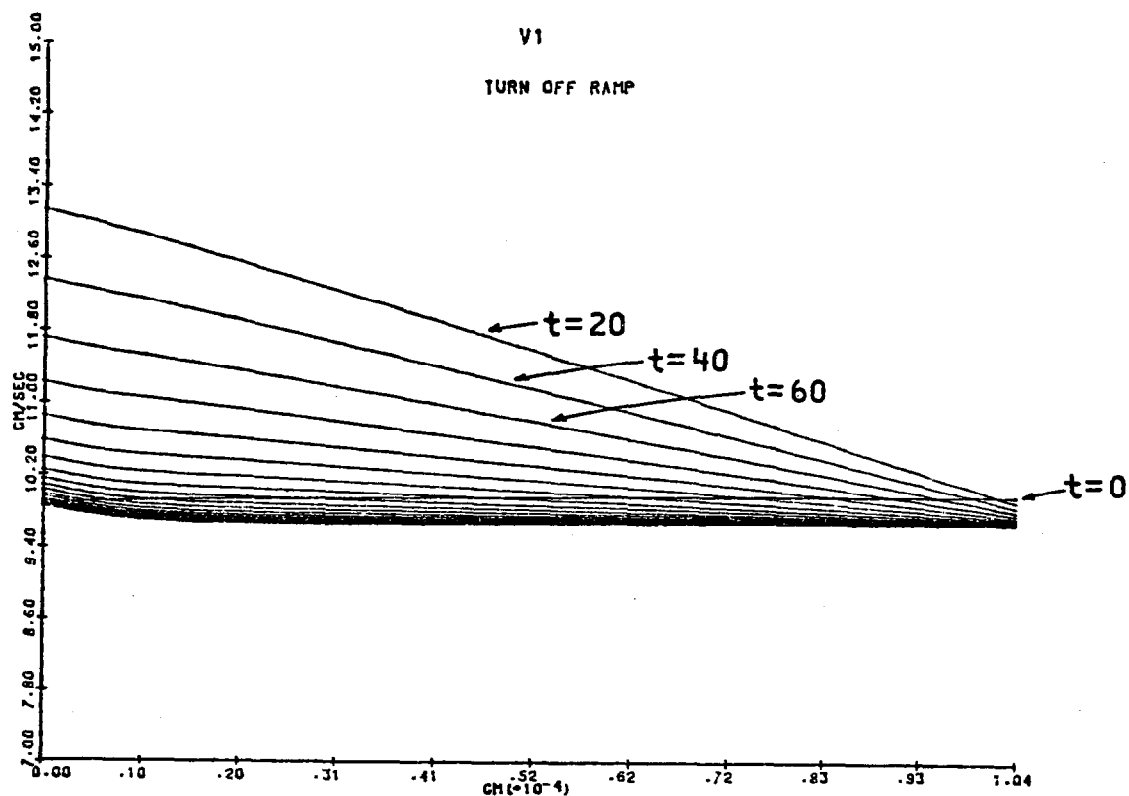


Figure 28. Velocity of single-phase stream as a function of x at 20 sec intervals.

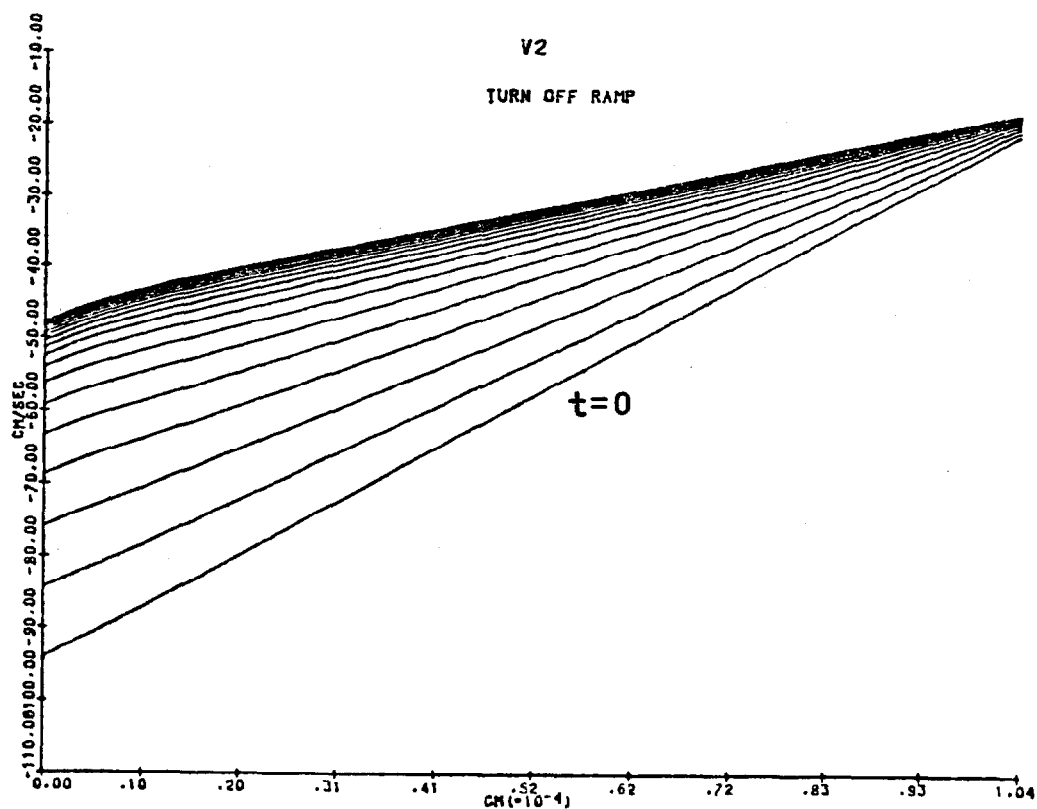


Figure 29. Velocity of two-phase stream as a function of x at 20 sec intervals.

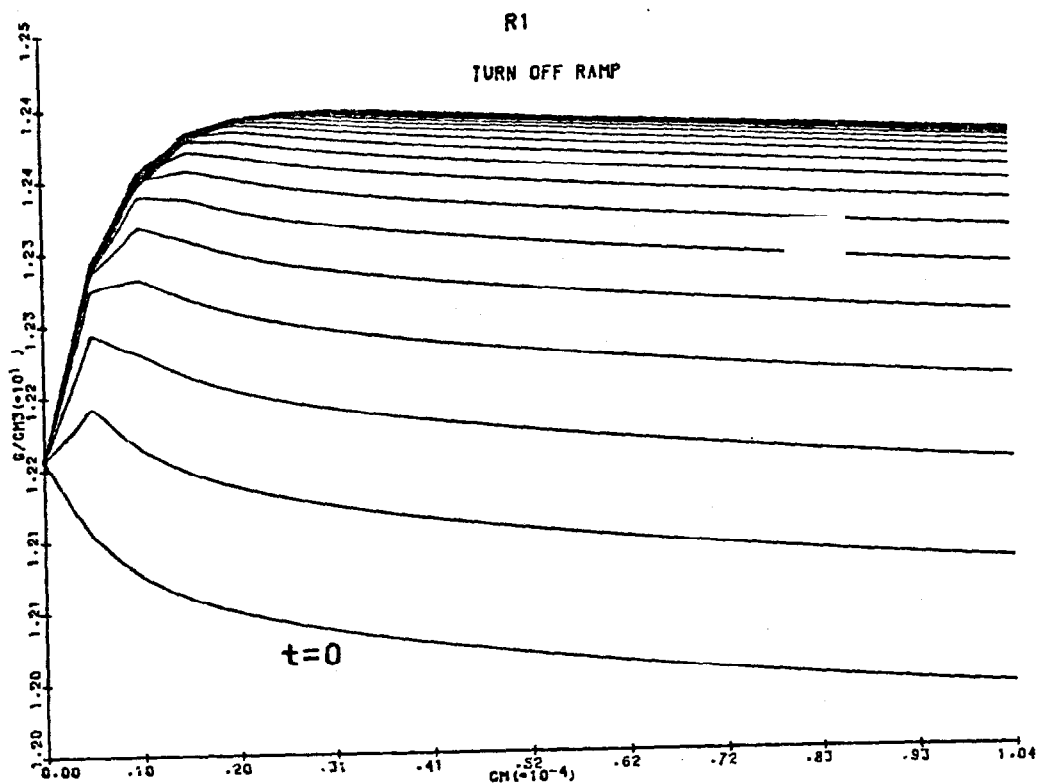


Figure 30. Density of single-phase stream as a function of x at 20 sec intervals.

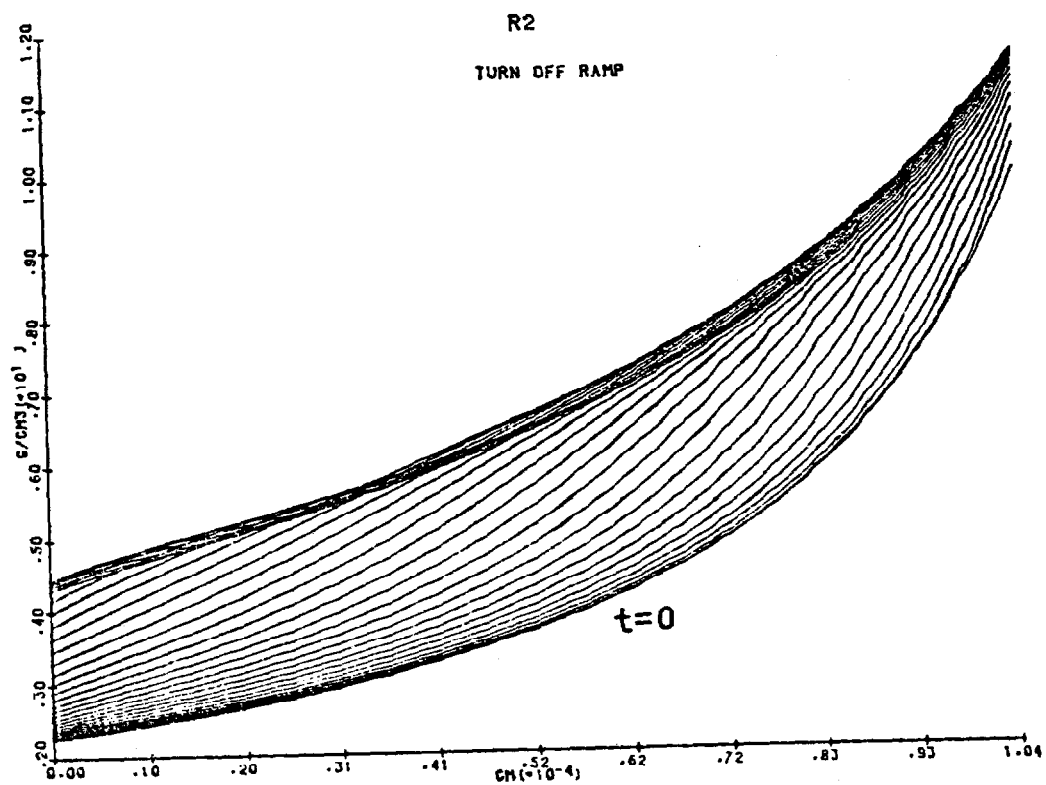


Figure 31. Density of two-phase stream as a function of x at 20 sec intervals.

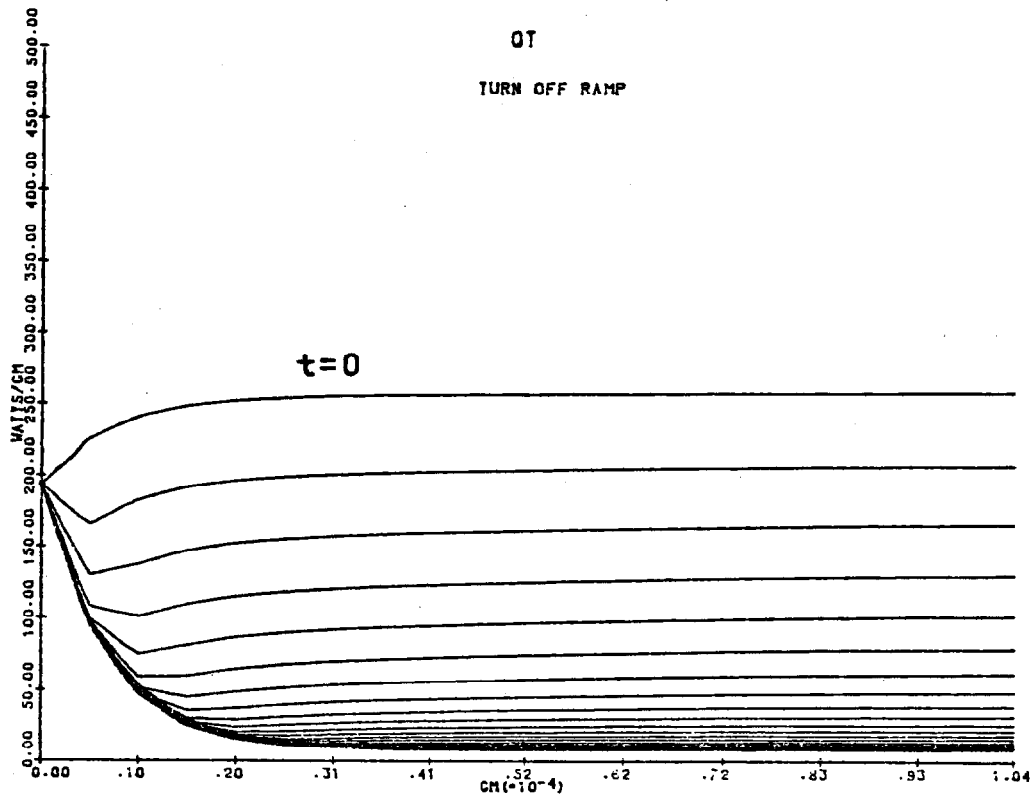


Figure 32. Heat transfer as a function of x at 20 sec intervals.

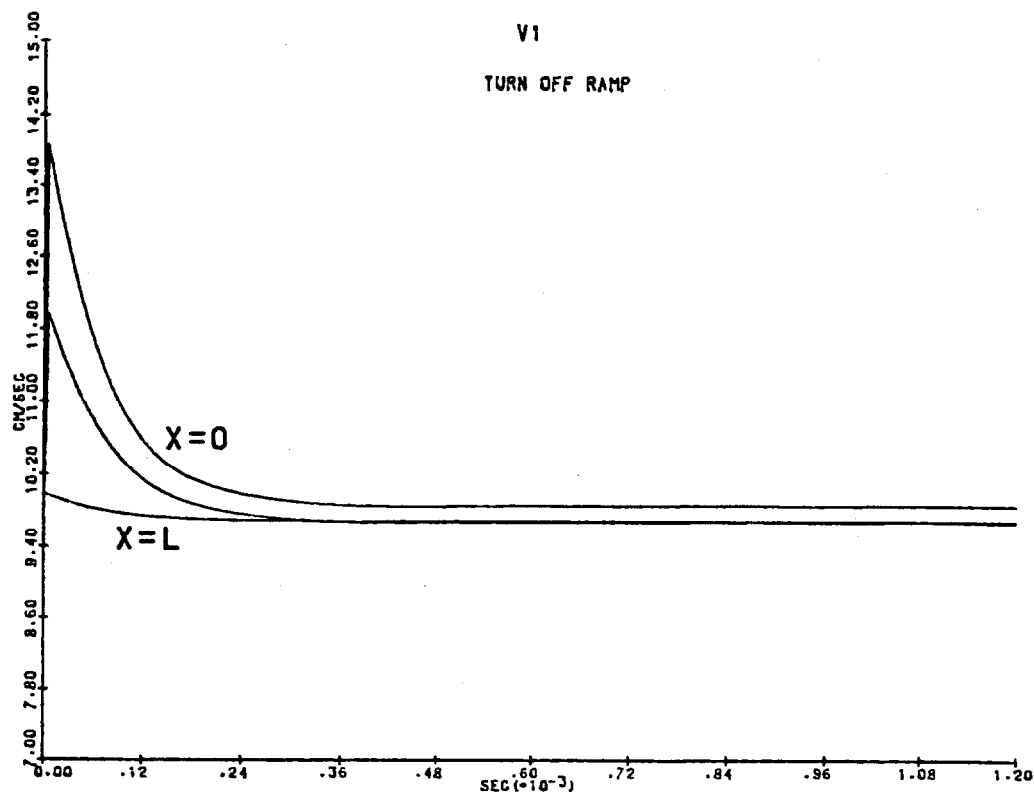


Figure 33. Velocity of single-phase stream as a function of time at $x=0$, $L/2$, and L .

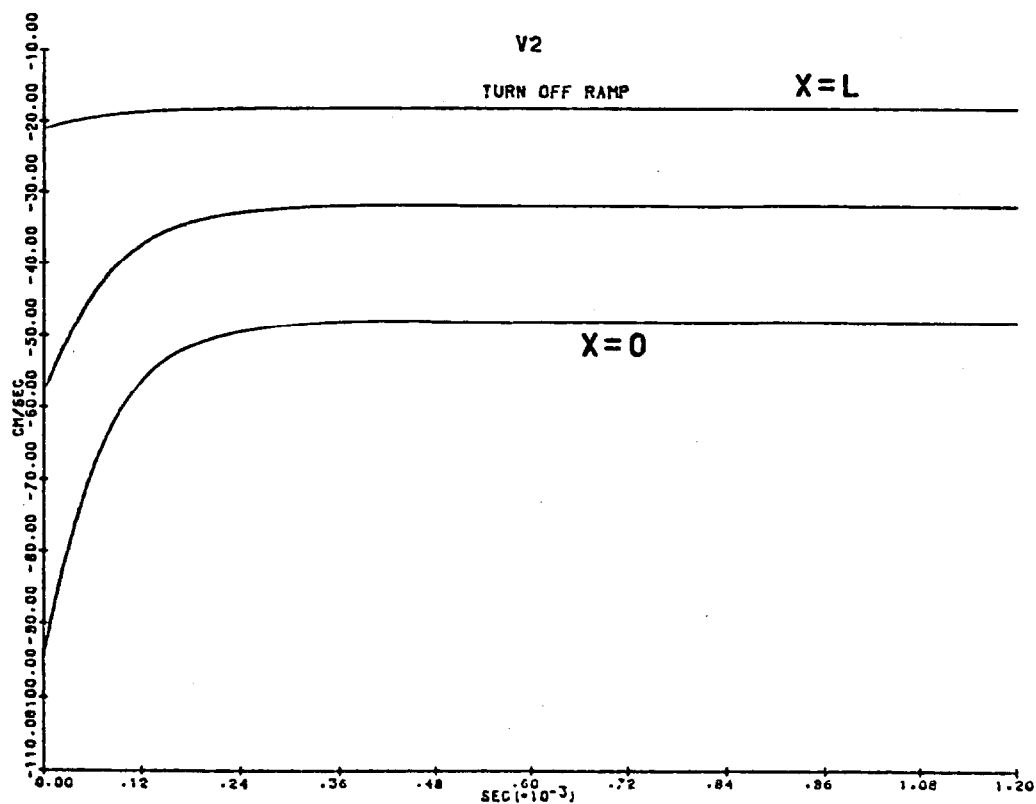


Figure 34. Velocity of two-phase stream as a function of time at $x=0$, $L/2$, and L .

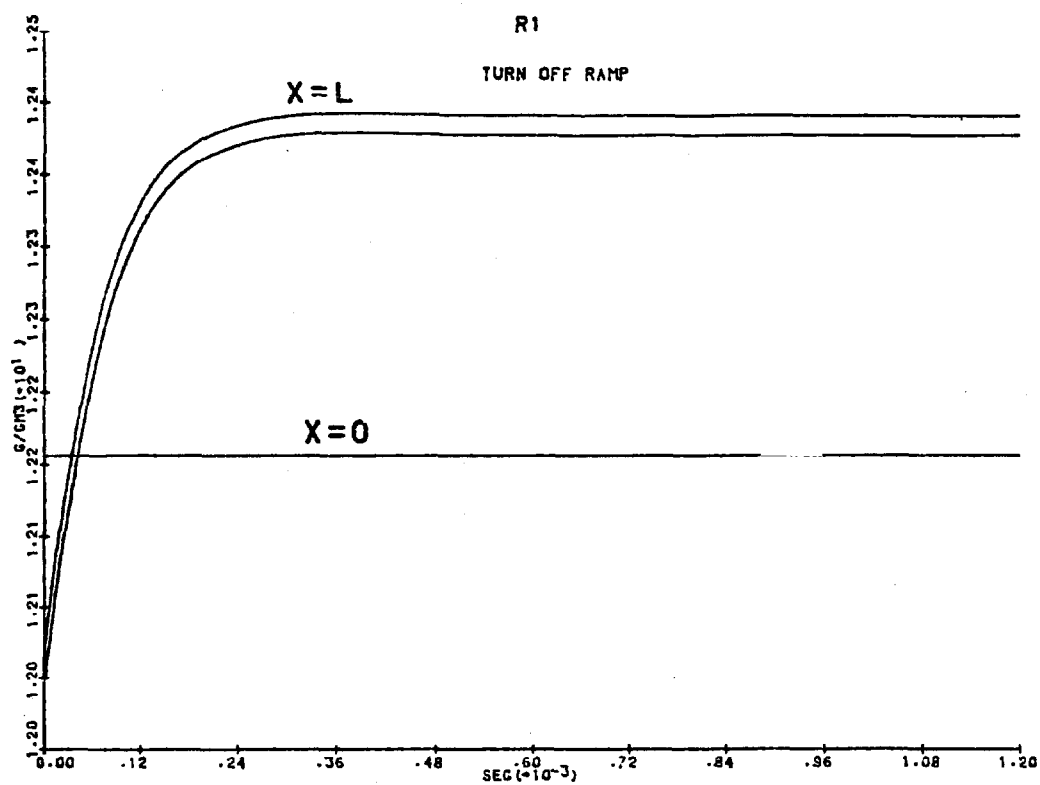


Figure 35. Density of single-phase stream as a function of time at $x=0$, $L/2$, and L .

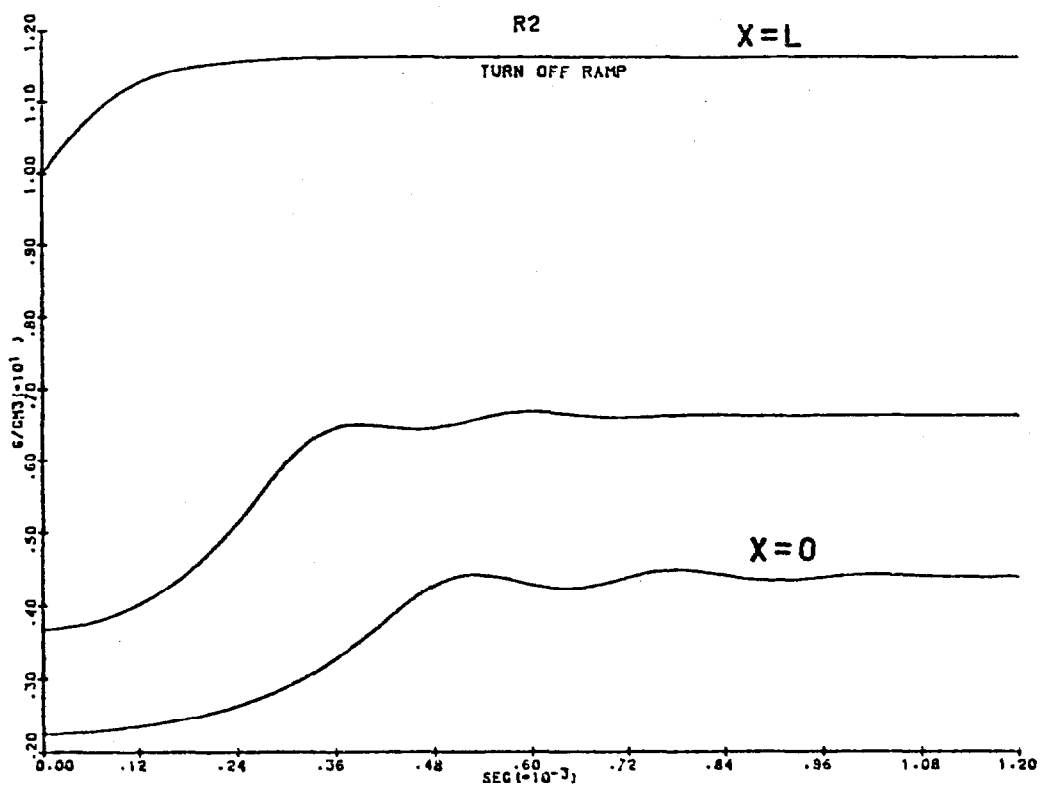


Figure 36. Density of two-phase stream as a function of time at $x=0$, $L/2$, and L .

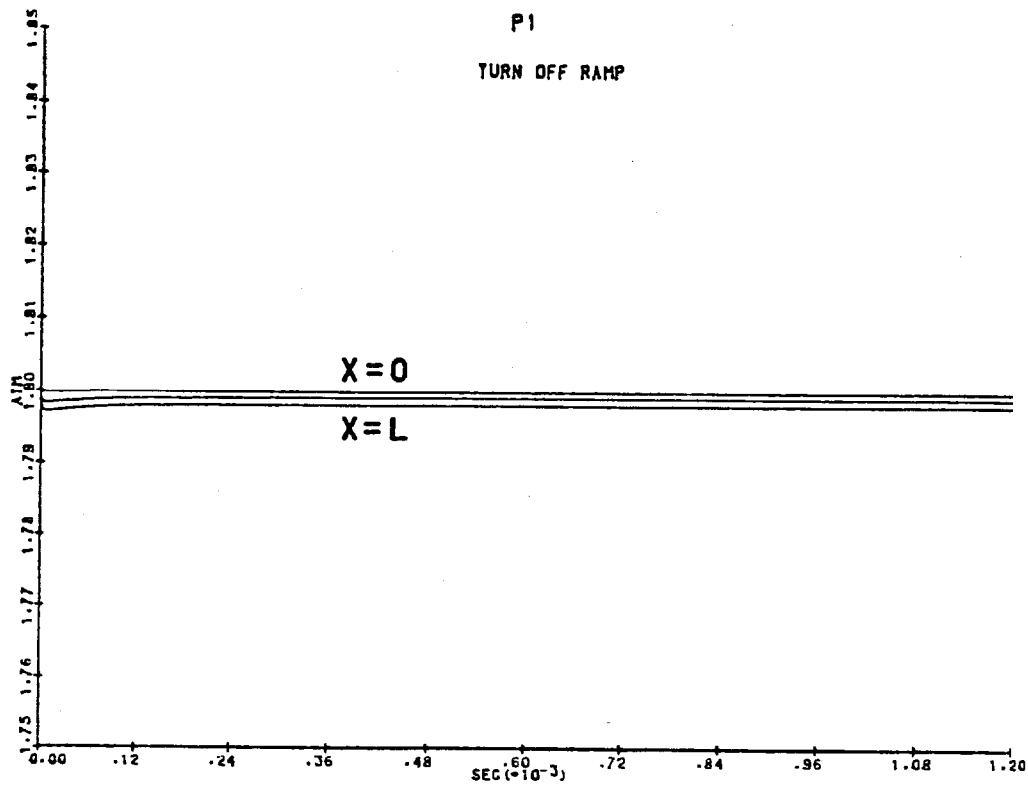


Figure 37. Pressure of single-phase stream as a function of time at $x=0$, $L/2$, and L .

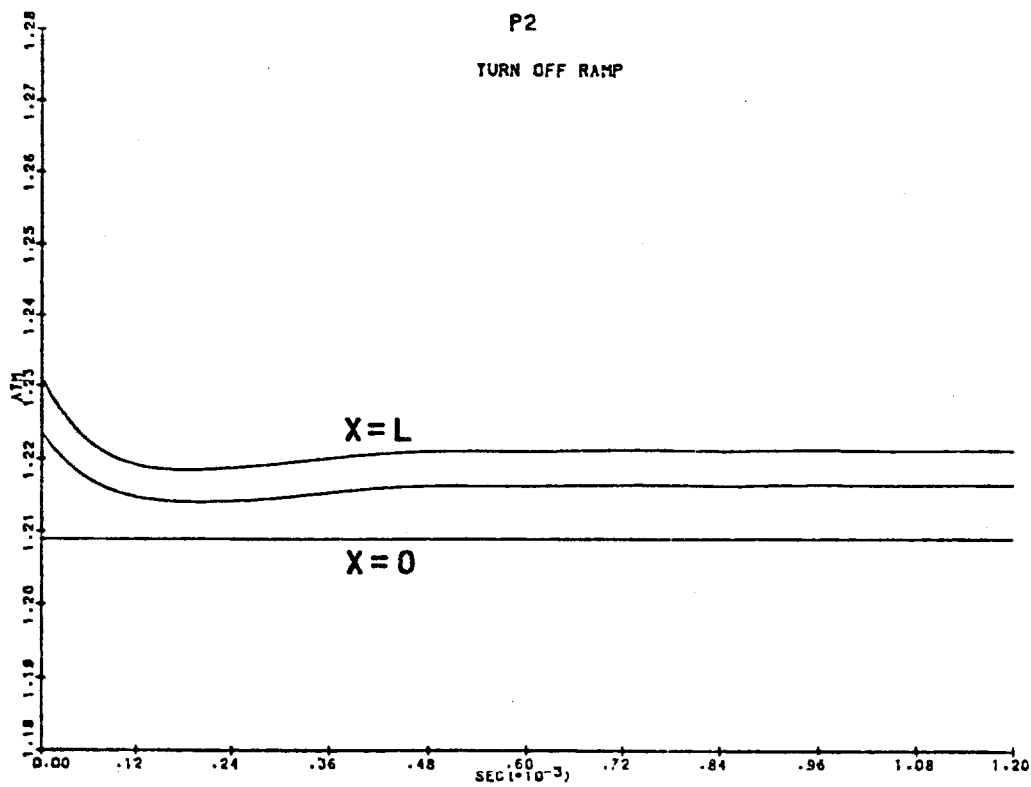


Figure 38. Pressure of two-phase stream as a function of time at $x=0$, $L/2$, and L .

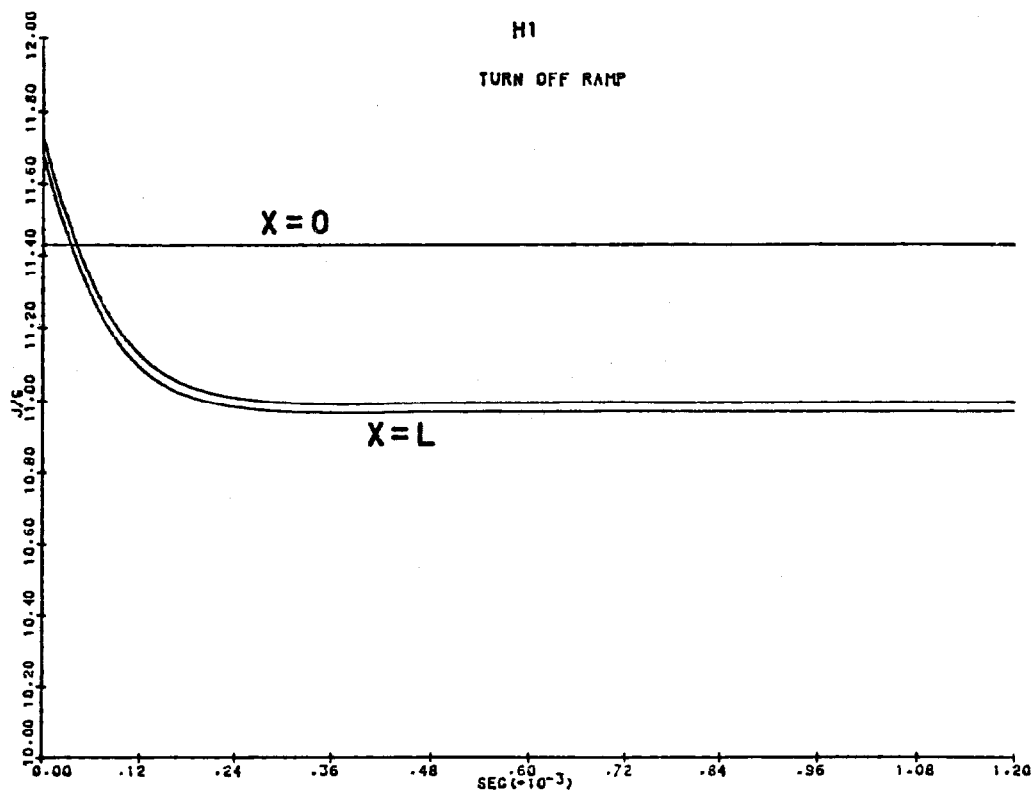


Figure 39. Enthalpy of single-phase stream as a function of time at $x=0$, $L/2$, and L .

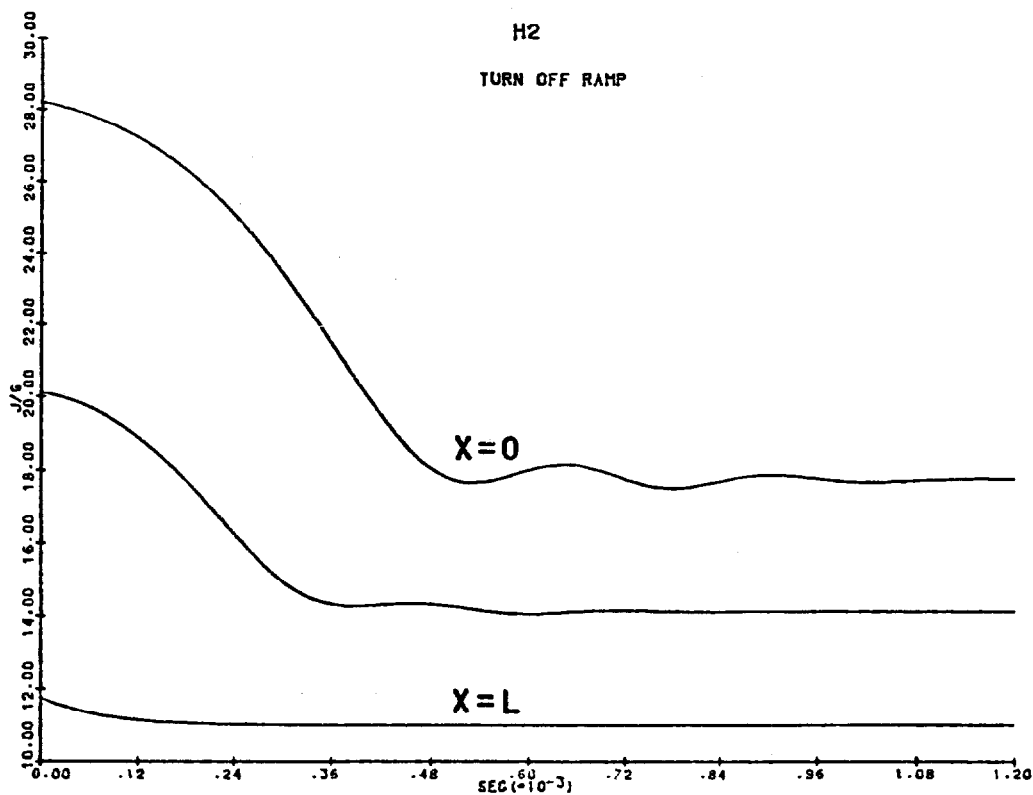


Figure 40. Enthalpy of two-phase stream as a function of time at $x=0$, $L/2$, and L .

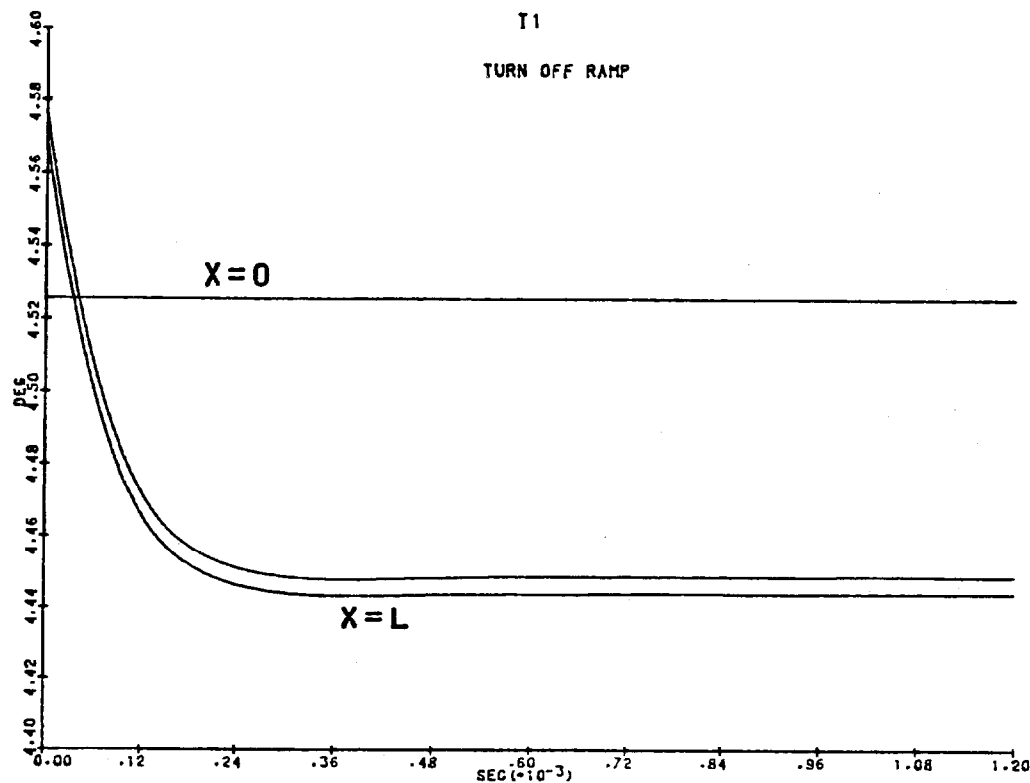


Figure 41. Temperature of single-phase stream as a function of time at $x=0$, $L/2$, and L .

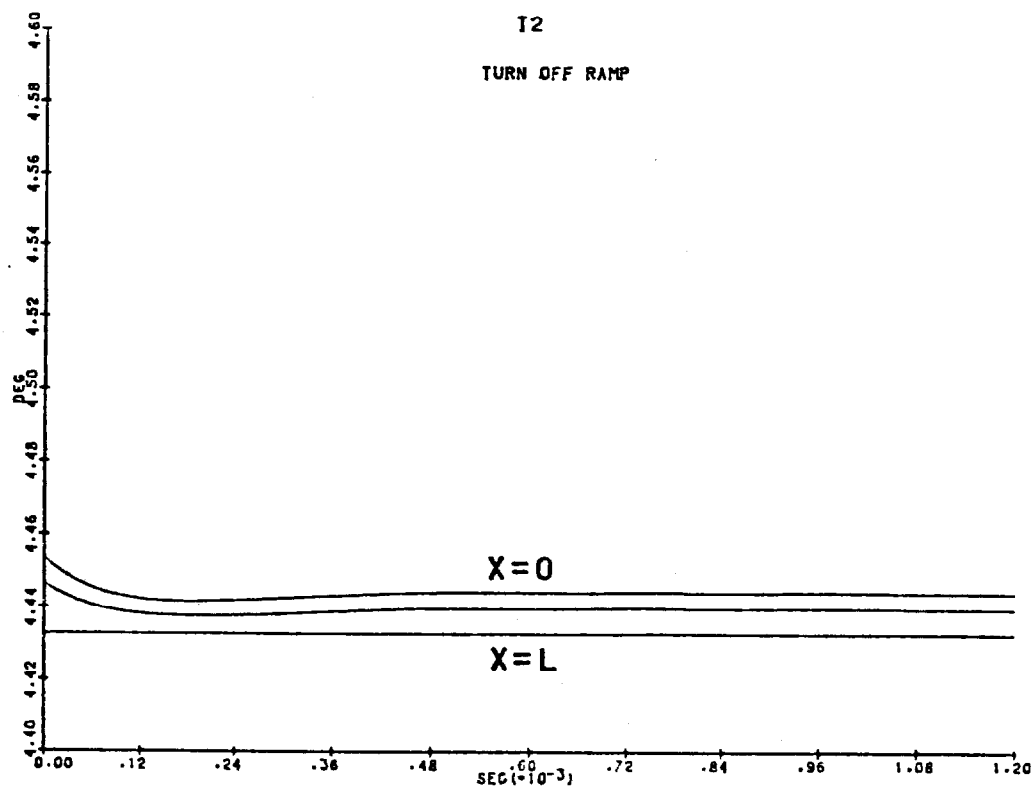


Figure 42. Temperature of two-phase stream as a function of time at $x=0$, $L/2$, and L .

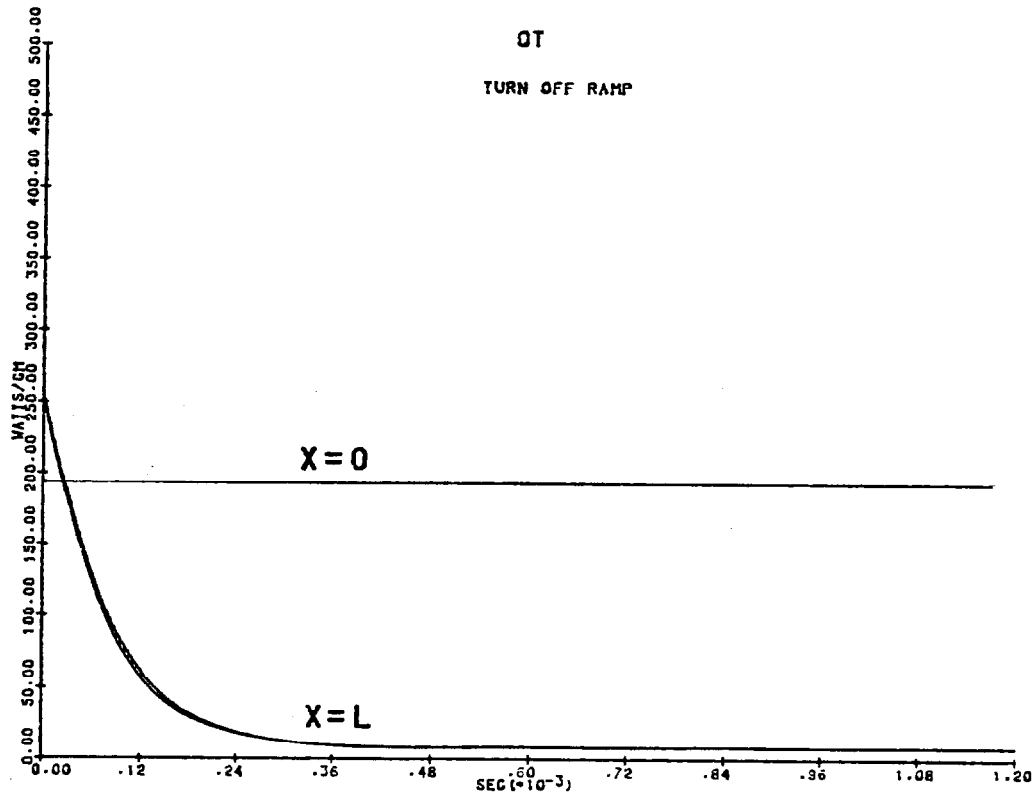


Figure 43. Heat transfer as a function of time at $x=0$, $L/2$, and L .